

**Numerical Methods of
Electromagnetic Field Theory I (NFT I)
Numerische Methoden der
Elektromagnetischen Feldtheorie I (NFT I) /**

3rd Lecture / 3. Vorlesung

Dr.-Ing. René Marklein

marklein@uni-kassel.de

<http://www.tet.e-technik.uni-kassel.de>

<http://www.uni-kassel.de/fb16/tet/marklein/index.html>

**Universität Kassel
Fachbereich Elektrotechnik / Informatik
(FB 16)
Fachgebiet Theoretische Elektrotechnik
(FG TET)
Wilhelmshöher Allee 71
Büro: Raum 2113 / 2115
D-34121 Kassel**

**University of Kassel
Dept. Electrical Engineering / Computer
Science (FB 16)
Electromagnetic Field Theory
(FG TET)
Wilhelmshöher Allee 71
Office: Room 2113 / 2115
D-34121 Kassel**

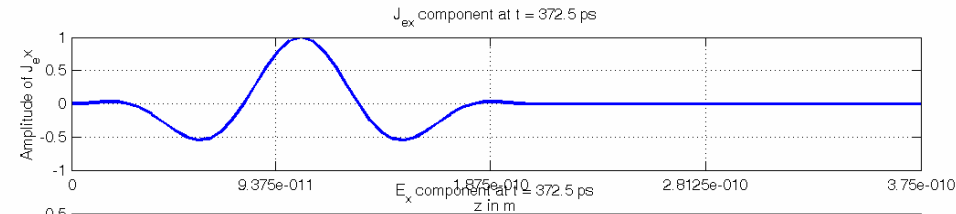
FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

Electric current density excitation: broadband pulse /
Elektrische Stromdichteanregung: breitbandiger Impuls

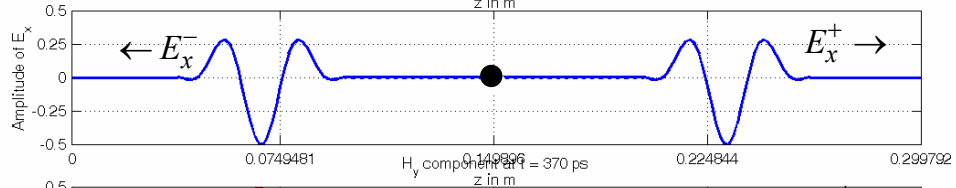
$$J_{ex}(z = z_0, t) \sim f_{RC2}(t) \rightarrow E_x(z, t) \sim f_{RC2} \left[t \mp \frac{z - z_0}{c_0} \right]$$

Snapshots / Schnappschüsse

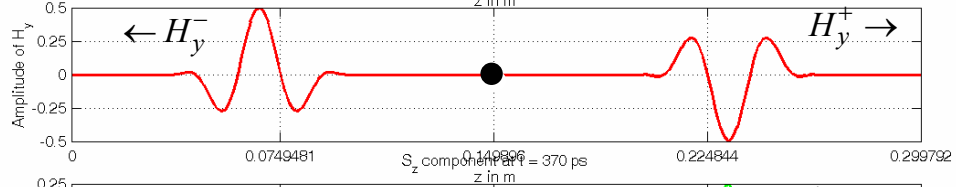
$f_{RC2}(t)$



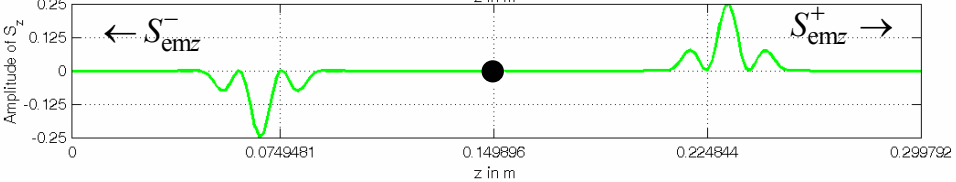
$\hat{E}_x(z, t_1)$



$\hat{H}_y(z, t_1)$



$\hat{S}_{emz}(z, t_1)$



● Source point / Quellpunkt

**Numerical Results – Validation /
Numerische Ergebnisse – Validierung**

Numerical Results / Numerische Ergebnisse



Validation / Validierung

Compare numerical results with analytical solutions or with other numerical solutions. / Vergleiche die numerischen Ergebnisse mit analytischen Lösungen oder anderen numerischen Lösungen

**Numerical Results – Validation /
Numerische Ergebnisse – Validierung**

**1. Plane Wave Solution of the Homogeneous Case –
No sources, no boundaries! /
Ebene Wellen als Lösung des homogenen Falles –
Keine Quellen, keine Ränder!**

Gives the correct characteristic, but not the correct amplitude and no reflections at the boundaries! / Gibt die korrekte Charakteristik, aber nicht die korrekte Amplitude und keine Reflexionen an den Rändern wieder!

**2. Green's Function Solution of the Inhomogeneous Case –
“Point” source, but no boundaries,
if we use the free-space Green's function! /
Lösung über Greensche Funktion für den inhomogenen Fall –
„Punkt“quelle, aber keine Ränder, wenn wir die
Greensche Funktion für den Freiraum verwenden!**

Gives the correct characteristic and correct amplitude, but no reflections at the boundaries! / Gibt die korrekte Charakteristik und die korrekte Amplitude, aber keine Reflexionen an den Rändern wieder!

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

Homogeneous scalar 1-D wave equation for the electric field strength / Homogene, skalare 1D-Wellengleichung für die elektrische Feldstärke

$$\frac{\partial^2}{\partial z^2} E_x(z,t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z,t) = 0$$

Splitting of the 1D wave operator /
Aufspaltung des 1D-Wellenoperators

$$\left(\frac{\partial^2}{\partial z^2} - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \right) E_x(z,t) = 0$$

$$\left(\frac{\partial}{\partial z} + \frac{1}{c_0} \frac{\partial}{\partial t} \right) \left(\frac{\partial}{\partial z} - \frac{1}{c_0} \frac{\partial}{\partial t} \right) E_x(z,t)$$

Hyperbolic partial differential equation /
Hyperbolische partielle Differentialgleichung

$$\left(\frac{\partial}{\partial z} + \frac{1}{c_0} \frac{\partial}{\partial t} \right) \left(\frac{\partial}{\partial z} - \frac{1}{c_0} \frac{\partial}{\partial t} \right) E_x(z,t) = 0$$

One-way wave equation /
"One-way" Wellengleichung

$$\left(\frac{\partial}{\partial z} + \frac{1}{c_0} \frac{\partial}{\partial t} \right) E_x(z,t) = 0$$

$$\left(\frac{\partial}{\partial z} - \frac{1}{c_0} \frac{\partial}{\partial t} \right) E_x(z,t) = 0$$

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

$$\begin{aligned}\left(\frac{\partial}{\partial z} + \frac{1}{c_0} \frac{\partial}{\partial t}\right) E_x(z, t) &= \left(\frac{\partial}{\partial z} + \frac{1}{c_0} \frac{\partial}{\partial t}\right) E_0^+ \left(z, t - \frac{z}{c_0}\right) \\ &= \frac{\partial}{\partial z} E_0^+ \left(z, t - \frac{z}{c_0}\right) + \frac{1}{c_0} \frac{\partial}{\partial t} E_0^+ \left(z, t - \frac{z}{c_0}\right) \\ &= -\frac{1}{c_0} E_0^+ \left(z, t - \frac{z}{c_0}\right) + \frac{1}{c_0} E_0^+ \left(z, t - \frac{z}{c_0}\right) \\ &= 0\end{aligned}$$

$$\begin{aligned}\left(\frac{\partial}{\partial z} - \frac{1}{c_0} \frac{\partial}{\partial t}\right) E_x(z, t) &= \left(\frac{\partial}{\partial z} - \frac{1}{c_0} \frac{\partial}{\partial t}\right) E_0^- \left(z, t + \frac{z}{c_0}\right) \\ &= \frac{\partial}{\partial z} E_0^- \left(z, t + \frac{z}{c_0}\right) - \frac{1}{c_0} \frac{\partial}{\partial t} E_0^- \left(z, t + \frac{z}{c_0}\right) \\ &= \frac{1}{c_0} E_0^- \left(z, t + \frac{z}{c_0}\right) - \frac{1}{c_0} E_0^- \left(z, t + \frac{z}{c_0}\right) \\ &= 0\end{aligned}$$

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

Homogeneous scalar 1-D wave equation for the electric field strength / Homogene, skalare 1D-Wellengleichung für die elektrische Feldstärke

$$\frac{\partial^2}{\partial z^2} E_x(z,t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z,t) = 0$$

$$\left(\frac{\partial}{\partial z} + \frac{1}{c_0} \frac{\partial}{\partial t} \right) \left(\frac{\partial}{\partial z} - \frac{1}{c_0} \frac{\partial}{\partial t} \right) E_x(z,t) = 0$$

Solution is a left and right propagating plane wave /
Lösung ist eine nach links und rechts laufende ebene Welle

$$E_x(z,t) = E_0 \left(z, t \mp \frac{z}{c_0} \right) \\ = \underbrace{E_0^+ \left(z, t - \frac{z}{c_0} \right)}_{\text{right-propagating}} + \underbrace{E_0^- \left(z, t + \frac{z}{c_0} \right)}_{\text{left-propagating}}$$

A wave, which propagates for increasing time t in positive z direction /
Eine Welle, die sich für zunehmende Zeit t in positive z -Richtung ausbreitet

A wave, which propagates for increasing time t in negative z direction /
Eine Welle, die sich für zunehmende Zeit t in negative z -Richtung ausbreitet

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

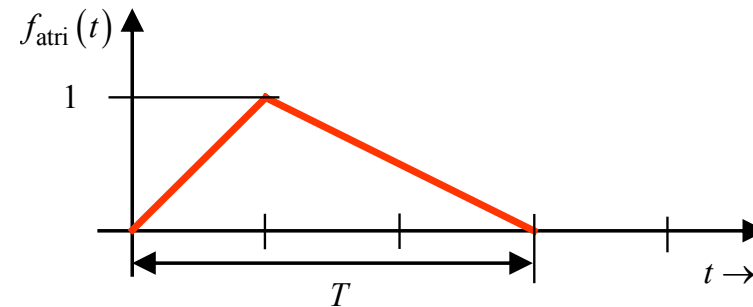
Consider an asymmetric triangular pulse /
Betrachte einen asymmetrischen Dreiecksimpuls

$$E_x(z = z_0, t) = E_0 f_{\text{atri}}(t)$$

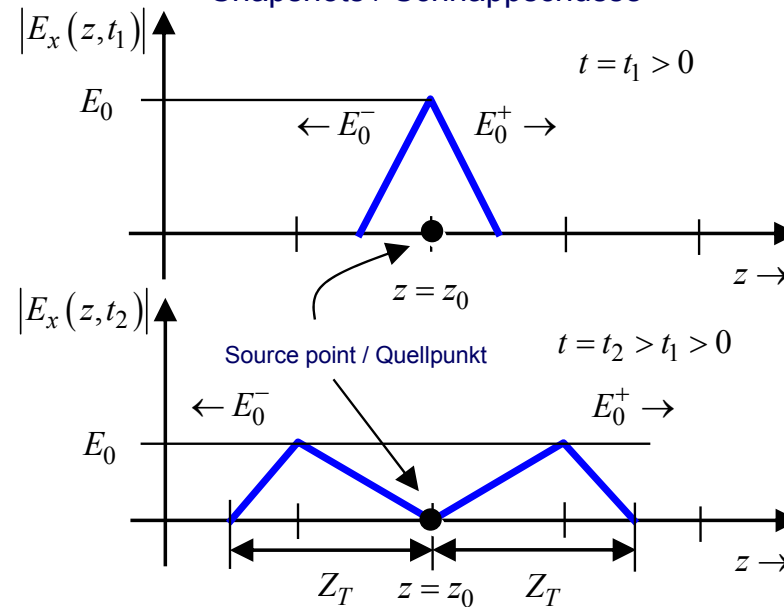
This means, that the solution for all z and t is given by / Dies bedeutet, dass die Lösung für alle z und t gegeben ist durch

$$\begin{aligned} E_x(z, t) &= E_0 \left(z, t \mp \frac{z - z_0}{c_0} \right) \\ &= E_0 f_{\text{atri}} \left(z, t \mp \frac{z - z_0}{c_0} \right) \\ &= E_0 f_{\text{atri}} \left(z, t - \frac{z - z_0}{c_0} \right) + E_0 f_{\text{atri}} \left(z, t + \frac{z - z_0}{c_0} \right) \\ &= E_0^+ \left(z, t - \frac{z - z_0}{c_0} \right) + E_0^- \left(z, t + \frac{z - z_0}{c_0} \right) \end{aligned}$$

Excitation function / Anregungsfunktion



Snapshots / Schnappschüsse



FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

$$E_x(z, t) = E_0 f_{\text{atri}} \left(z, t \mp \frac{z - z_0}{c_0} \right)$$

$$t \mp \frac{z - z_0}{c_0} = 0$$

$$t = \pm \frac{z - z_0}{c_0}$$

$$c_0 t = \pm z \mp z_0$$

$$\pm c_0 t = z - z_0$$

$$z(t) = z_0 \pm c_0 t$$

$$z^+(t) = z_0 + c_0 t$$

$$z^-(t) = z_0 - c_0 t$$

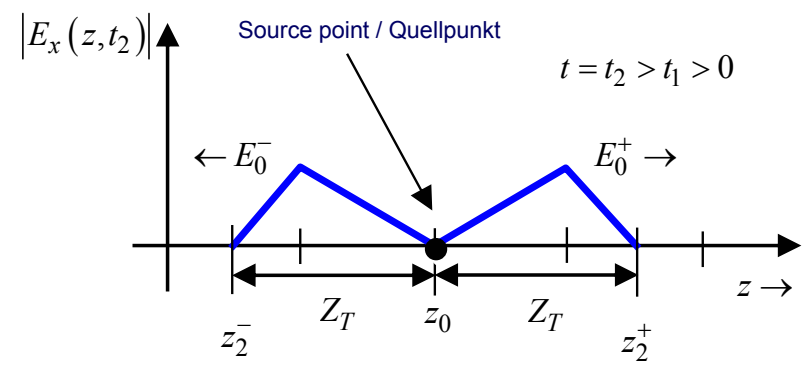
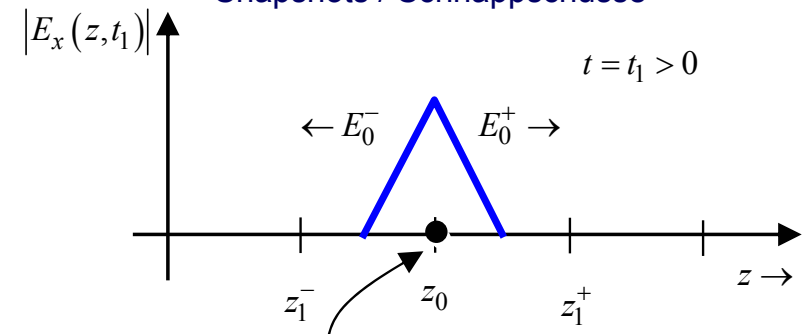
$$z_1^+(t_1) = z_0 + c_0 t_1$$

$$z_1^-(t_1) = z_0 - c_0 t_1$$

$$z_2^+(t_2) = z_0 + c_0 t_2$$

$$z_2^-(t_2) = z_0 - c_0 t_2$$

Snapshots / Schnappschüsse



FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

$$\frac{\partial^2}{\partial z^2} E_x(z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z, t) = 0$$

$$E_x(z, t) = E_0 \left(z, t \mp \frac{z}{c_0} \right)$$

$$\frac{\partial}{\partial t} H_y(z, t) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z, t)$$

?

$$S_{\text{emz}}(z, t) = E_x(z, t) H_y(z, t)$$

FD Method – 1-D Helmholtz Equation (Reduced Wave Equation) FD-Methode – 1D Helmholtz-Gleichung (Schwingungsgleichung)

Homogeneous scalar 1-D wave equation /
 Homogene, skalare 1D-Wellengleichung

$$\frac{\partial^2}{\partial z^2} E_x(z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z, t) = 0$$

1-D Fourier transform with
 regard to time t /
 1D Fourier-Transformation
 bezüglich der Zeit t

$$E_x(z, \omega) = \int_{t=-\infty}^{\infty} E_x(z, t) e^{j\omega t} dt$$

$$= FT_t \{ E_x(z, t) \}$$

$$E_x(z, \omega) \bullet \dashv \circ E_x(z, t)$$

1-D inverse Fourier transform with
 regard to circular frequency ω /
 1D inverse Fourier-Transformation
 bezüglich der Kreisfrequenz ω

$$E_x(z, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} E_x(z, \omega) e^{-j\omega t} d\omega$$

$$= FT_{\omega}^{-1} \{ E_x(z, \omega) \}$$

$$E_x(z, t) \circ \bullet E_x(z, \omega)$$

$$E_x(z, \omega) \bullet \dashv \circ E_x(z, t)$$

$$-j\omega \bullet \dashv \circ \frac{\partial}{\partial t}$$

$$-\omega^2 \bullet \dashv \circ \frac{\partial^2}{\partial t^2}$$

FD Method – 1-D Helmholtz Equation (Reduced Wave Equation) FD-Methode – 1D Helmholtz-Gleichung (Schwingungsgleichung)

Homogeneous scalar 1-D wave equation /
 Homogene, skalare 1D-Wellengleichung

$$\frac{\partial^2}{\partial z^2} E_x(z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z, t) = 0$$

Solution in the time domain /
 Lösung im Zeitbereich

$$E_x(z, t) = E_0 \left(z, t \mp \frac{z}{c_0} \right)$$

Homogeneous scalar 1-D Helmholtz wave
 equation (reduced wave equation) /
 Homogene, skalare 1D Helmholtz-Gleichung
 (Schwingungsgleichung)

$$\frac{\partial^2}{\partial z^2} E_x(z, \omega) - \frac{1}{c_0^2} (-\omega^2) E_x(z, \omega) = 0$$

$$\frac{\partial^2}{\partial z^2} E_x(z, \omega) + \underbrace{\frac{\omega^2}{c_0^2}}_{=k_0^2} E_x(z, \omega) = 0$$

$$\frac{\partial^2}{\partial z^2} E_x(z, \omega) + k_0^2 E_x(z, \omega) = 0$$

Solution in the frequency domain /
 Lösung im Frequenzbereich

$$E_x(z, \omega) = E_0(\omega) e^{\pm j k_0 z}$$

$$E_x(z, t) \longleftrightarrow E_x(z, \omega)$$

$$\frac{\partial^2}{\partial t^2} \longleftrightarrow -\omega^2$$



FD Method – 1-D Helmholtz Equation (Reduced Wave Equation) FD-Methode – 1D Helmholtz-Gleichung (Schwingungsgleichung)

Maxwell's equations in the time domain / Maxwell'sche Gleichungen im Zeitbereich

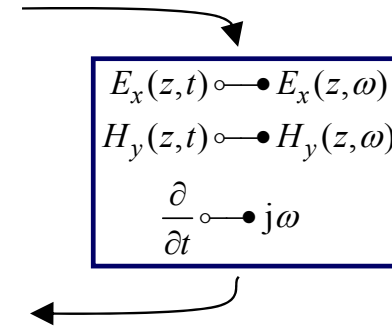
$$\frac{\partial}{\partial t} H_y(z, t) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z, t)$$

$$\frac{\partial}{\partial t} E_x(z, t) = -\frac{1}{\varepsilon_0} \frac{\partial}{\partial z} H_y(z, t)$$

Maxwell's equations in the frequency domain / Maxwell'sche Gleichungen im Frequenzbereich

$$-j\omega H_y(z, \omega) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z, \omega)$$

$$-j\omega E_x(z, \omega) = -\frac{1}{\varepsilon_0} \frac{\partial}{\partial z} H_y(z, \omega)$$



Electric field strength: plane wave /
Elektrische Feldstärke: ebene Welle

$$E_x(z, \omega) = E_0(\omega) e^{\pm jk_0 z}$$

$$H_y(z, \omega) = \frac{1}{j\omega\mu_0} \frac{\partial}{\partial z} E_x(z, \omega) = \frac{1}{j\omega\mu_0} \frac{\partial}{\partial z} E_0(\omega) e^{\pm jk_0 z} = \frac{1}{j\omega\mu_0} E_0(\omega) \underbrace{\frac{\partial}{\partial z} e^{\pm jk_0 z}}_{=\pm jk_0 e^{\pm jk_0 z}} = \pm \frac{jk_0}{j\omega\mu_0} E_0(\omega) e^{\pm jk_0 z}$$

$$\omega / c_0 = \omega \sqrt{\varepsilon_0 \mu_0}$$

$$= \pm \frac{\overbrace{k_0}}{\omega\mu_0} E_0(\omega) e^{\pm jk_0 z} = \pm \frac{1}{Z_0} E_0(\omega) e^{\pm jk_0 z}$$

Magnetic field strength: plane wave /
Magnetische Feldstärke: ebene Welle

FD Method – 1-D Helmholtz Equation (Reduced Wave Equation) FD-Methode – 1D Helmholtz-Gleichung (Schwingungsgleichung)

Homogeneous scalar 1-D wave equation in the time domain / Homogene, skalare 1D-Wellengleichung im Zeitbereich

$$\frac{\partial^2}{\partial z^2} E_x(z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z, t) = 0$$

$$\frac{\partial^2}{\partial z^2} H_y(z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(z, t) = 0$$

Solution of the 1-D wave equation in the time domain / Lösung der homogenen 1D-Wellengleichung im Zeitbereich

$$E_x(z, t) = E_0 \left(z, t \mp \frac{z}{c_0} \right)$$

$$H_y(z, t) = H_0 \left(z, t \mp \frac{z}{c_0} \right)$$

Homogeneous, scalar 1-D Helmholtz equation in the frequency domain / Homogene, skalare 1D-Helmholtz-Gleichung im Frequenzbereich

$$\frac{\partial^2}{\partial z^2} E_x(z, \omega) + k_0^2 E_x(z, \omega) = 0$$

$$\frac{\partial^2}{\partial z^2} H_y(z, \omega) + k_0^2 H_y(z, \omega) = 0$$

Solution of the 1-D Helmholtz equation in the frequency domain / Lösung der homogenen 1D-Helmholtz-Gleichung im Frequenzbereich

$$E_x(z, \omega) = E_0(\omega) e^{\pm j k_0 z}$$

$$H_y(z, \omega) = H_0(\omega) e^{\pm j k_0 z}$$

$$= \pm \frac{1}{Z_0} E_0(\omega) e^{\pm j k_0 z}$$

Solution of the 1-D wave equation for the magnetic field strength in terms of the electric field strength / Lösung der homogenen 1D-Wellengleichung für die magnetische Feldstärke als Funktion der elektrischen Feldstärke

$$H_y(z, t) = \pm \frac{1}{Z_0} E_0 \left(z, t \mp \frac{z}{c_0} \right)$$

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

Homogeneous scalar 1-D wave equations /
Homogene, skalare 1D-Wellengleichungen

$$\frac{\partial^2}{\partial z^2} E_x(z,t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(z,t) = 0$$

$$\frac{\partial^2}{\partial z^2} H_y(z,t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(z,t) = 0$$

Solutions / Lösungen

$$E_x(z,t) = E_0 \left(z, t \mp \frac{z}{c_0} \right)$$

$$H_y(z,t) = \pm \frac{1}{Z_0} E_0 \left(z, t \mp \frac{z}{c_0} \right)$$

Poynting vector / Poynting-Vector

$$S_{\text{emz}}(z,t) = E_x(z,t)H_y(z,t)$$

$$= E_0 \left(z, t \mp \frac{z}{c_0} \right) H_0 \left(z, t \mp \frac{z}{c_0} \right)$$

$$= E_0 \left(z, t \mp \frac{z}{c_0} \right) \left[\pm \frac{1}{Z_0} E_0 \left(z, t \mp \frac{z}{c_0} \right) \right]$$

$$= \pm \frac{1}{Z_0} E_0^2 \left(z, t \mp \frac{z}{c_0} \right)$$

$$S_{\text{emz}}(z,t) = \pm \frac{1}{Z_0} E_0^2 \left(z, t \mp \frac{z}{c_0} \right)$$

$$= \underbrace{\frac{1}{Z_0} E_0^2 \left(z, t - \frac{z}{c_0} \right)}_{S_{\text{emz}}^+(z,t)} - \underbrace{\frac{1}{Z_0} E_0^2 \left(z, t + \frac{z}{c_0} \right)}_{S_{\text{emz}}^-(z,t)}$$

$$= S_{\text{emz}}^+(z,t) + S_{\text{emz}}^-(z,t)$$

Poynting vector of the two plane waves /
Poynting-Vektor der beiden ebenen Wellen

$$S_{\text{emz}}^+(z,t) = \frac{1}{Z_0} E_0^2 \left(z, t - \frac{z}{c_0} \right)$$

$$S_{\text{emz}}^-(z,t) = -\frac{1}{Z_0} E_0^2 \left(z, t + \frac{z}{c_0} \right)$$

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

$$E_x(z, t) = E_0^+ \left(z, t - \frac{z - z_0}{c_0} \right) + E_0^- \left(z, t + \frac{z - z_0}{c_0} \right)$$

$$f_{RC2}(t)$$

$$H_y(z, t) = \left[\frac{E_0^+ \left(z, t - \frac{z - z_0}{c_0} \right)}{Z_0} \right] + \left[-\frac{E_0^- \left(z, t + \frac{z - z_0}{c_0} \right)}{Z_0} \right]$$

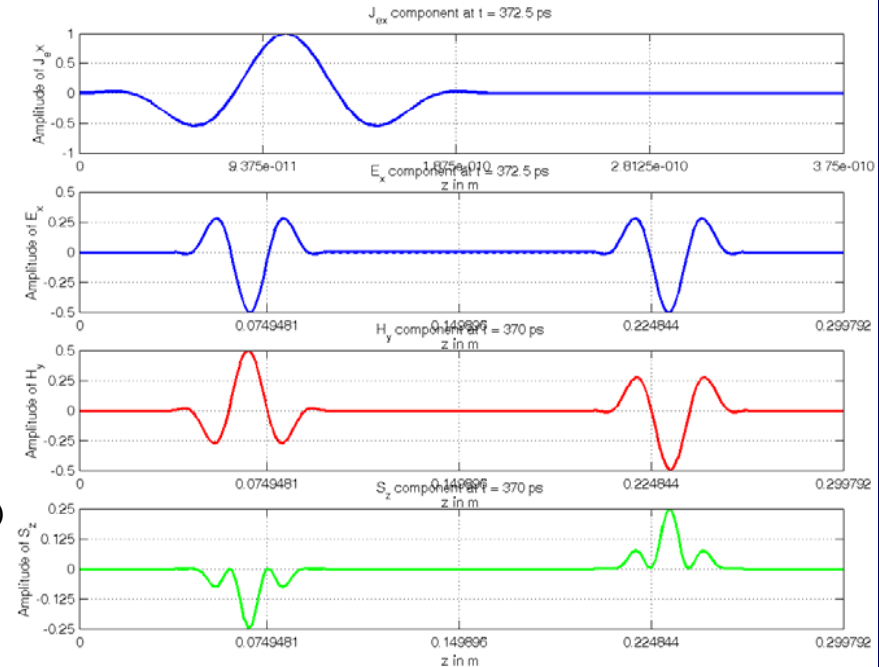
$$\hat{E}_x(z, t_1)$$

$$\hat{H}_y(z, t_1)$$

$$S_{\text{emz}}(z, t) = S_{\text{emz}}^+(z, t) + S_{\text{emz}}^-(z, t)$$

$$= \left[\frac{E_0^2 \left(z, t - \frac{z - z_0}{c_0} \right)}{Z_0} \right] + \left[-\frac{E_0^2 \left(z, t + \frac{z - z_0}{c_0} \right)}{Z_0} \right]$$

$$\hat{S}_{\text{emz}}(z, t_1)$$



The plane wave solution gives the correct characteristic of the wave field, but the amplitude is not correct! This means we can not verify the numerical results with the plane wave solution of the homogeneous wave equation, because the simulated problem correspond to the solution of the inhomogeneous wave equation. /

Die Ebene-Wellen-Lösung gibt die korrekte Charakteristik des Wellenfeldes wieder, aber die Amplitude der Wellenanteile ist nicht korrekt! Dies bedeutet, dass man die numerischen Resultate mit der Ebenen-Wellen-Lösung nicht vollständig verifizieren kann, da die simulierte Situation mit der Lösung der inhomogenen Wellengleichung korrespondiert.

Electromagnetic Field of a Point Source Excitation in 1-D / Elektromagnetisches Feld einer Punktquellenanregung in 1D

Integral representation /
Integraldarstellung

$$E_x(z, \omega) = j\omega\mu_0 \int_{z'=-\infty}^{\infty} G(z-z', \omega) J_{ex}(z', \omega) dz'$$

1-D scalar Green's function in the frequency domain /
1D skalare Greensche Funktion im Frequenzbereich

$$\begin{aligned} G(z-z', \omega) &= \frac{1}{2} \left[\text{PV} \frac{j}{k_0} + \pi c_0 \delta(z-z') \right] e^{jk_0|z-z'|} \\ &= \frac{1}{2} \left[\text{PV} \frac{j}{k_0} e^{jk_0|z-z'|} + \pi c_0 \delta(z-z') \right] \\ &= \frac{1}{2} \left[-\text{PV} \frac{c_0}{j\omega} e^{jk_0|z-z'|} + \pi c_0 \delta(z-z') \right] \\ &= \frac{c_0}{2} \left[-\text{PV} \frac{1}{j\omega} e^{jk_0|z-z'|} + \pi \delta(z-z') \right] \end{aligned}$$

1-D scalar Green's function
in the time domain /
1D skalare Greensche Funktion
im Zeitbereich

$$G(z, t) = \frac{c_0}{2} u\left(t - \frac{|z|}{c_0}\right)$$

Unit step function /
Einheitssprungfunktion

$$u(t) = \begin{cases} 0 & t < 0 \\ 1 & t > 0 \end{cases}$$

Electric current density /
Elektrische Stromdichte

$$\begin{aligned} J_{ex}(z, \omega) &= \delta(z-z_0) K_{ex}(z, \omega) \\ &= \delta(z-z_0) K_{ex}(z_0, \omega) \end{aligned}$$

Electric surface current density /
Elektrische Flächenstromdichte

$$K_{ex}(z_0, \omega)$$

Property of the delta-distribution /
Eigenschaft der Delta-Distribution

$$\delta(z-z_0) f(z) = \delta(z-z_0) f(z_0)$$

EM Field of a Point Source Excitation in 1-D / EM-Feld einer Punktquellenanregung in 1D

$$\begin{aligned}
 E_x(z, \omega) &= j\omega\mu_0 \int_{z'=-\infty}^{\infty} G(z-z', \omega) J_{\text{ex}}(z', \omega) dz' \\
 &= j\omega\mu_0 \int_{z'=-\infty}^{\infty} G(z-z', \omega) \delta(z'-z_0) K_{\text{ex}}(z', \omega) dz' \\
 &= j\omega\mu_0 G(z-z_0, \omega) K_{\text{ex}}(z_0, \omega)
 \end{aligned}$$

$$\begin{aligned}
 E_x(z, \omega) &\bullet \longrightarrow E_x(z, t) \\
 j\omega &\bullet \longrightarrow -\frac{\partial}{\partial t} \\
 G(z-z_0, \omega) &\bullet \longrightarrow G(z-z_0, t) \\
 K_{\text{ex}}(z_0, \omega) &\bullet \longrightarrow K_{\text{ex}}(z_0, t) \\
 G(z-z_0, \omega)K_{\text{ex}}(z_0, \omega) &\bullet \longrightarrow G(z-z_0, t) *_t K_{\text{ex}}(z_0, t)
 \end{aligned}$$

The asterisk “*_t” denotes convolution in time / Der Stern “*_t” bezeichnet eine Faltung in der Zeit

$$\begin{aligned}
 E_x(z, t) &= -\mu_0 \frac{\partial}{\partial t} \int_{t'=-\infty}^{\infty} G(z-z_0, t-t') K_{\text{ex}}(z_0, t') dt' \\
 &= -\frac{c_0\mu_0}{2} \frac{\partial}{\partial t} \int_{t'=-\infty}^{\infty} u\left(t-t'-\frac{|z-z_0|}{c_0}\right) K_{\text{ex}}(z_0, t') dt' \\
 &= -\frac{c_0\mu_0}{2} \int_{t'=-\infty}^{\infty} \left[\frac{\partial}{\partial t} u\left(t-t'-\frac{|z-z_0|}{c_0}\right) \right] K_{\text{ex}}(z_0, t') dt'
 \end{aligned}$$

EM Field of a Point Source Excitation in 1-D / EM-Feld einer Punktquellenanregung in 1D

$$E_x(z, t) = -\frac{c_0 \mu_0}{2} \int_{t'=-\infty}^{\infty} \left[\frac{\partial}{\partial t} u \left(t - t' - \frac{|z - z_0|}{c_0} \right) \right] K_{\text{ex}}(z_0, t') dt'$$

$$\frac{\partial}{\partial t} u \left(t - t' - \frac{|z - z_0|}{c_0} \right) = \delta \left(t - t' - \frac{|z - z_0|}{c_0} \right)$$

$$\begin{aligned} E_x(z, t) &= -\frac{c_0 \mu_0}{2} \int_{t'=-\infty}^{\infty} \delta \left(t - t' - \frac{|z - z_0|}{c_0} \right) K_{\text{ex}}(z_0, t') dt' \\ &= -\frac{c_0 \mu_0}{2} \int_{t'=-\infty}^{\infty} \delta \left(t' - \left(t - \frac{|z - z_0|}{c_0} \right) \right) K_{\text{ex}}(z_0, t') dt' \\ &= -\frac{c_0 \mu_0}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \end{aligned}$$

$$c_0 \mu_0 = Z_0 \approx 377 \, \Omega \quad \text{Wave impedance of free space (vacuum) /
Wellenwiderstand des Freiraumes (Vakuum)}$$

**Solution for the x component of the electric field strength /
Lösung für die x -Komponente der elektrischen Feldstärke**

$$E_x(z, t) = -\frac{Z_0}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right)$$

EM Field of a Point Source Excitation in 1-D / EM-Feld einer Punktquellenanregung in 1D

$$\begin{aligned}
 H_y(z, \omega) &= \frac{1}{j\omega\mu_0} \frac{\partial}{\partial z} E_x(z, \omega) \\
 &= \frac{1}{j\omega\mu_0} \frac{\partial}{\partial z} [j\omega\mu_0 G(z - z_0, \omega) K_{\text{ex}}(z_0, \omega)] \\
 &= \frac{\partial}{\partial z} G(z - z_0, \omega) K_{\text{ex}}(z_0, \omega)
 \end{aligned}$$



$$E_x(z, \omega) = j\omega\mu_0 G(z - z_0, \omega) K_{\text{ex}}(z_0, \omega)$$

$$\begin{aligned}
 \frac{\partial}{\partial t} u(t) &= \delta(t) \\
 \frac{\partial}{\partial t} u(-t) &= -\delta(t)
 \end{aligned}$$

$$\begin{aligned}
 H_y(z, t) &= \frac{\partial}{\partial z} \int_{t'=-\infty}^{\infty} G(z - z_0, t - t') K_{\text{ex}}(z_0, t') dt' \\
 &= \frac{c_0}{2} \frac{\partial}{\partial z} \int_{t'=-\infty}^{\infty} u\left(t - t' - \frac{|z - z_0|}{c_0}\right) K_{\text{ex}}(z_0, t') dt' \\
 &= \frac{c_0}{2} \int_{t'=-\infty}^{\infty} \left[\frac{\partial}{\partial z} u\left(t - t' - \frac{|z - z_0|}{c_0}\right) \right] K_{\text{ex}}(z_0, t') dt' \\
 &= -\frac{c_0}{2} \frac{\text{sgn}(z - z_0)}{c_0} \int_{t'=-\infty}^{\infty} \delta\left(t - t' - \frac{|z - z_0|}{c_0}\right) K_{\text{ex}}(z_0, t') dt' \\
 &= -\text{sgn}(z - z_0) \frac{1}{2} K_{\text{ex}}\left(z_0, t - \frac{|z - z_0|}{c_0}\right)
 \end{aligned}$$

$$\begin{aligned}
 &\frac{\partial}{\partial z} u\left(t - t' - \frac{|z - z_0|}{c_0}\right) \\
 &= \frac{\partial}{\partial z} u\left(t - t' - \frac{\text{sgn}(z - z_0)(z - z_0)}{c_0}\right) \\
 &= -\frac{\text{sgn}(z - z_0)}{c_0} \delta\left(t - t' - \frac{\text{sgn}(z - z_0)(z - z_0)}{c_0}\right) \\
 &= -\frac{\text{sgn}(z - z_0)}{c_0} \delta\left(t - t' - \frac{|z - z_0|}{c_0}\right)
 \end{aligned}$$

EM Field of a Point Source Excitation in 1-D / EM-Feld einer Punktquellenanregung in 1D

Solution for the y component of the magnetic field strength /
Lösung für die y -Komponente der magnetische Feldstärke

$$H_y(z, t) = -\frac{\text{sgn}(z - z_0)}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right)$$

Solution for the x component of the electric field strength /
Lösung für die x -Komponente der elektrischen Feldstärke

$$E_x(z, t) = -\frac{Z_0}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right)$$

Solution for the z component of the Poynting vector /
Lösung für die z -Komponente des Poynting-Vektors

$$\begin{aligned} S_{\text{emz}}(z, t) &= E_x(z, t)H_y(z, t) \\ &= \left[-\frac{Z_0}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \right] \left[-\frac{\text{sgn}(z - z_0)}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \right] \\ &= \text{sgn}(z - z_0) \frac{Z_0}{4} K_{\text{ex}}^2 \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \end{aligned}$$

EM Field of a Point Source Excitation in 1-D / EM-Feld einer Punktquellenanregung in 1D

Normalization of the field components / Normierung der Feldkomponenten

$$\begin{aligned}
 \Delta t &= \Delta t_{\text{ref}} \widehat{\Delta t} & \Delta t_{\text{ref}} &= \frac{\Delta x_{\text{ref}}}{c_{\text{ref}}} & \Delta z &= \Delta x_{\text{ref}} \widehat{\Delta z} & c &= c_{\text{ref}} \widehat{c} & \varepsilon &= \varepsilon_{\text{ref}} \widehat{\varepsilon} & \mu &= \mu_{\text{ref}} \widehat{\mu} & \mu_{\text{ref}} &= \mu_0 \\
 E_x &= E_{\text{ref}} \widehat{E}_x \\
 H_y &= H_{\text{ref}} \widehat{H}_y & H_{\text{ref}} &= \frac{E_{\text{ref}}}{c_{\text{ref}} \mu_{\text{ref}}} = \frac{\sqrt{\varepsilon_{\text{ref}} \mu_{\text{ref}}}}{\mu_{\text{ref}}} E_{\text{ref}} = \sqrt{\frac{\varepsilon_{\text{ref}}}{\mu_{\text{ref}}}} E_{\text{ref}} = \frac{E_{\text{ref}}}{Z_{\text{ref}}} \\
 S_{\text{em}z} &= S_{\text{em ref}} \widehat{S}_{\text{em}z} & S_{\text{em ref}} &= E_{\text{ref}} H_{\text{ref}} = \frac{E_{\text{ref}}^2}{Z_{\text{ref}}} & c_{\text{ref}} &= c_0 \\
 J_{\text{ex}} &= J_{\text{e ref}} \widehat{J}_{\text{ex}} & J_{\text{e ref}} &= \frac{\varepsilon_{\text{ref}}}{\Delta t_{\text{ref}}} E_{\text{ref}} & \Delta x_{\text{ref}} &= \Delta z \\
 \delta(z) &= \frac{1}{\Delta x_{\text{ref}}} \widehat{\delta}(z) & \varepsilon_{\text{ref}} &= \varepsilon_0 \\
 K_{\text{ex}} &= K_{\text{e ref}} \widehat{K}_{\text{ex}} & K_{\text{e ref}} &= \Delta x_{\text{ref}} J_{\text{e ref}} & \mu_{\text{ref}} &= \mu_0 \\
 & & & & Z_{\text{ref}} &= Z_0 \\
 & & & & K_{\text{e ref}} &= 1 \text{ A/m}
 \end{aligned}$$

EM Field components / EM-Feldkomponenten

$$\begin{aligned}
 K_{\text{ex}}(z_0, t) &= K_{\text{e ref}} f_{RC2}(t) \\
 E_x^\pm(z, t) &= -\frac{Z_0}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \\
 H_y^\pm(z, t) &= \mp \frac{1}{2} K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \\
 S_{\text{em}z}^\pm(z, t) &= \pm \frac{Z_0}{4} \left[K_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \right]^2
 \end{aligned}$$

Normalized EM field components / Normierte EM-Feldkomponenten

$$\begin{aligned}
 \widehat{K}_{\text{ex}}(z_0, t) &= f_{RC2}(t) \\
 \widehat{E}_x^\pm(z, t) &= -\frac{1}{2} \widehat{K}_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \\
 \widehat{H}_y^\pm(z, t) &= \mp \frac{1}{2} \widehat{K}_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \\
 \widehat{S}_{\text{em}z}^\pm(z, t) &= \pm \frac{1}{4} \left[\widehat{K}_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \right]^2
 \end{aligned}$$

FD Method – 1-D Wave Equation – Example / FD-Methode – 1D Wellengleichung – Beispiel

$$\widehat{K}_{\text{ex}}(z_0, t) = f_{\text{RC2}}(t)$$

$$\widehat{E}_x^\pm(z, t) = -\frac{1}{2} \widehat{K}_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right)$$

$$\widehat{H}_y^\pm(z, t) = \mp \frac{1}{2} \widehat{K}_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right)$$

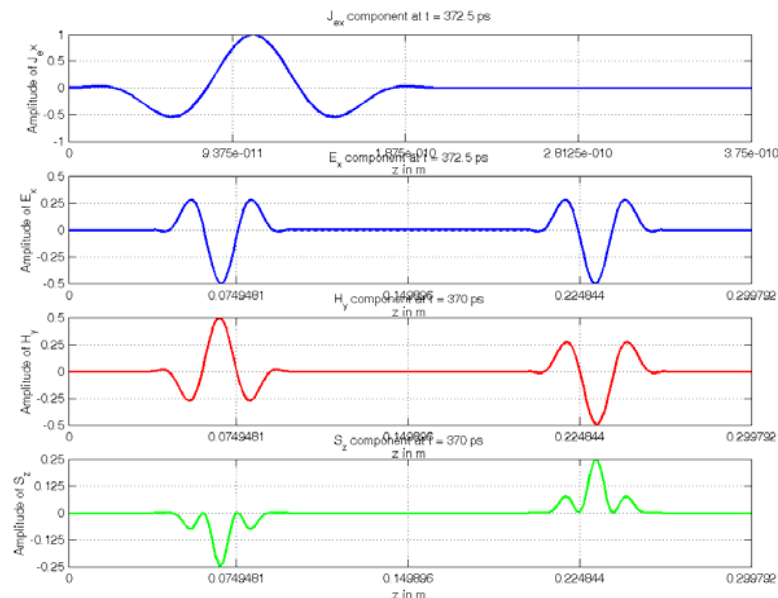
$$\widehat{S}_{\text{emz}}^\pm(z, t) = \pm \frac{1}{4} \left[\widehat{K}_{\text{ex}} \left(z_0, t - \frac{|z - z_0|}{c_0} \right) \right]^2$$

$$f_{\text{RC2}}(t)$$

$$\widehat{E}_x(z, t_1)$$

$$\widehat{H}_y(z, t_1)$$

$$\widehat{S}_{\text{emz}}(z, t_1)$$



The Green's function method gives the solution of the 1-D simulation area excited by a "point" source, which is in 1-D a singular electric surface current source. The singular source is independent of x and y . The reference solution gives the correct characteristic and correct amplitudes. But the solution doesn't account for the reflections at the boundaries, because we used the free-space Green's function. /

Die Methode der Greenschen Funktion ermöglicht die Lösung des vorliegenden Problems, der Anregung des 1D-Simulationsgebietes durch eine „Punkt“quelle, die genauer gesagt in 1D eine singuläre elektrische Flächenstromdichte ist. Da die singuläre Quelle von x und y unabhängig ist. Die Charakteristik und Amplitude stimmt überein, nur die Reflexionen an den Rändern fehlen, was an der Verwendung der Greenschen Funktion für den Freiraum liegt.

**End of Lecture 3 /
Ende der 3. Vorlesung**