

**Numerical Methods of  
Electromagnetic Field Theory I (NFT I)  
Numerische Methoden der  
Elektromagnetischen Feldtheorie I (NFT I) /**

**4th Lecture / 4. Vorlesung**

**Dr.-Ing. René Marklein**

[marklein@uni-kassel.de](mailto:marklein@uni-kassel.de)

<http://www.tet.e-technik.uni-kassel.de>

<http://www.uni-kassel.de/fb16/tet/marklein/index.html>

**Universität Kassel  
Fachbereich Elektrotechnik / Informatik  
(FB 16)  
Fachgebiet Theoretische Elektrotechnik  
(FG TET)  
Wilhelmshöher Allee 71  
Büro: Raum 2113 / 2115  
D-34121 Kassel**

**University of Kassel  
Dept. Electrical Engineering / Computer  
Science (FB 16)  
Electromagnetic Field Theory  
(FG TET)  
Wilhelmshöher Allee 71  
Office: Room 2113 / 2115  
D-34121 Kassel**

## FD Method – Properties / FD-Methode - Eigenschaften

✚ Spatial and Temporal Discretization /  
Räumliche und zeitliche Diskretisierung

$$\Delta z = ?$$
$$\Delta t = ?$$

✚ Consistency /  
Konsistenz

✚ Dissipation /  
Dissipation

✚ Stability Condition /  
Stabilitätsbedingung

$$\Delta t = f(\Delta z)$$

✚ Convergence /  
Konvergenz

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

### Stability by the *von Neumann's* method (Fourier series method):

Insert a complex monofrequent (monochromatic) plane wave into the discrete FD equations and analyze the spectral radius of the amplification matrix, where the spectral radius must be smaller equal one.

### Stabilität durch die *von Neumannsche* Methode (Fourier-Reihen-Methode):

Setze eine komplex monofrequente (monochromatische) ebene Welle in die diskreten FD-Gleichungen ein und analysiere den spektralen Radius der Verstärkungsmatrix, wobei der spektrale Radius kleinergleich Eins sein muss.

Complex monofrequent (monochromatic) plane wave /  
Komplex monofrequente (monochromatische) ebene Welle

$$E_x(\mathbf{R}, t) = E_0(\omega_0, \hat{\mathbf{k}}) e^{-j(\omega_0 t - \mathbf{k} \cdot \mathbf{R})}$$

$$= E_0(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 t} e^{j\mathbf{k} \cdot \mathbf{R}}$$

$$\{\mathbf{W}\}^{(n+1)} = [\mathbf{G}]_{1D}^{\text{FD}} \{\mathbf{W}\}^{(n)} \quad [\mathbf{G}]_{1D}^{\text{FD}} : \text{Amplification matrix / Verstärkungsmatrix}$$

Spectral radius /  
Spektraler Radius  $\rho([\mathbf{G}]_{1D}^{\text{FD}}) \leq 1$  of the matrix /  
der Matrix  $[\mathbf{G}]_{1D}^{\text{FD}}$

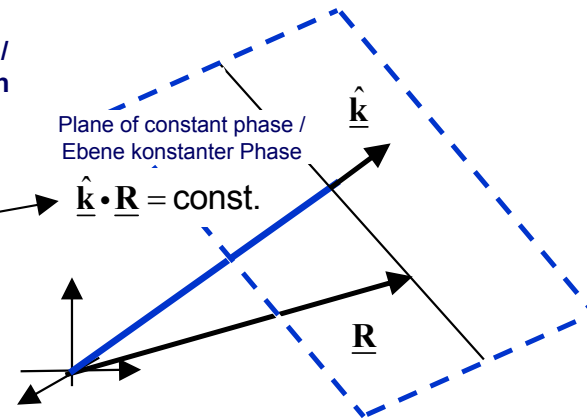
where /  
wobei  $\rho([\mathbf{G}]_{1D}^{\text{FD}}) = \max_{n=1, \dots, N} |v_n([\mathbf{G}]_{1D}^{\text{FD}})|$   $v_n([\mathbf{G}]_{1D}^{\text{FD}})$ :  $n$ th eigenvalue of the matrix  $[\mathbf{G}]_{1D}^{\text{FD}}$   
 $n$ -ter Eigenwert der Matrix  $[\mathbf{G}]_{1D}^{\text{FD}}$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

Monofrequent (monochromatic) plane wave in the time domain /  
Monofrequente (monochromatische) ebene Welle im Zeitbereich

$$E_x(\underline{\mathbf{R}}, t) = E_0(\omega_0, \hat{\mathbf{k}}) e^{-j(\omega_0 t - k \hat{\mathbf{k}} \cdot \underline{\mathbf{R}})}$$

$$= E_0(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 t} e^{jk \hat{\mathbf{k}} \cdot \underline{\mathbf{R}}}$$



Wave vector /  
Wellenvektor

$$\underline{\mathbf{k}} = k_x \underline{\mathbf{e}}_x + k_y \underline{\mathbf{e}}_y + k_z \underline{\mathbf{e}}_z = k_z \underline{\mathbf{e}}_z$$

Magnitude of the wave vector /  
Betrag des Wellenvektors

$$|\underline{\mathbf{k}}| = \sqrt{\underline{\mathbf{k}} \cdot \underline{\mathbf{k}}} = \sqrt{k_x^2 + k_y^2 + k_z^2} = \sqrt{k_z^2} = |k_z| = k$$

Wavenumber /  
Wellenzahl

$$k = \frac{\omega_0}{c}$$

Circular frequency /  
Kreisfrequenz

$$\omega_0 = 2\pi f_0$$

Propagation direction /  
Ausbreitungsrichtung

$$\hat{\mathbf{k}} = \frac{\underline{\mathbf{k}}}{|\underline{\mathbf{k}}|} = \frac{k_z \underline{\mathbf{e}}_z}{|k_z|} = \frac{\text{sgn}(k_z) |k_z| \underline{\mathbf{e}}_z}{k} = \frac{\text{sgn}(k_z) k \underline{\mathbf{e}}_z}{k} = \text{sgn}(k_z) \underline{\mathbf{e}}_z$$

Phase of the plane wave /  
Phase der ebenen Welle

$$k \hat{\mathbf{k}} \cdot \underline{\mathbf{R}} = k \text{sgn}(k_z) \underline{\mathbf{e}}_z \cdot (x \underline{\mathbf{e}}_x + y \underline{\mathbf{e}}_y + z \underline{\mathbf{e}}_z) = k \text{sgn}(k_z) z \underline{\mathbf{e}}_z \cdot \underline{\mathbf{e}}_z = k \text{sgn}(k_z) z$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

Insert discrete plane wave / Setze die diskrete ebene Welle

$$\hat{E}_x^{(n_z, n_t)} = \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \underbrace{e^{jk n_z \Delta z}}_{=\exp(n_z)} e^{-j\omega_0 n_t \Delta t} = \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{-j\omega_0 n_t \Delta t}$$

into the FD scheme / in das FD-Schema ein

$$\hat{E}_x^{(n_z, n_t+1)} = 2\hat{E}_x^{(n_z, n_t)} - \hat{E}_x^{(n_z, n_t-1)} + (\widehat{\Delta t})^2 \left[ \hat{E}_x^{(n_z+1, n_t)} - 2\hat{E}_x^{(n_z, n_t)} + \hat{E}_x^{(n_z-1, n_t)} \right]$$

with / mit

$$\begin{aligned} \hat{E}_x^{(n_z, n_t+1)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \underbrace{e^{jk n_z \Delta z}}_{=\exp(n_z)} e^{-j\omega_0 (n_t+1) \Delta t} \\ &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{-j\omega_0 (n_t+1) \Delta t} & \hat{E}_x^{(n_z+1, n_t)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{jk \Delta z} e^{-j\omega_0 n_t \Delta t} \\ \hat{E}_x^{(n_z, n_t)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{-j\omega_0 n_t \Delta t} & \hat{E}_x^{(n_z, n_t)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{-j\omega_0 n_t \Delta t} \\ \hat{E}_x^{(n_z, n_t-1)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{j\omega_0 (n_t-1) \Delta t} & \hat{E}_x^{(n_z-1, n_t)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) \exp(n_z) e^{-jk \Delta z} e^{-j\omega_0 n_t \Delta t} \end{aligned}$$

it follows / folgt

$$\begin{aligned} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 (n_z+1) \Delta t} &= 2\hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 (n_z-1) \Delta t} \\ &\quad + (\widehat{\Delta t})^2 \left[ \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{jk \Delta z} - 2\hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) + \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-jk \Delta z} \right] e^{-j\omega_0 n_z \Delta t} \\ &= 2 \left[ 1 - (\widehat{\Delta t})^2 \right] \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 (n_z-1) \Delta t} \\ &\quad + (\widehat{\Delta t})^2 \left[ \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{jk \Delta z} + \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-jk \Delta z} \right] e^{-j\omega_0 n_z \Delta t} \end{aligned}$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

$$\begin{aligned}
 \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z+1)\Delta t} &= 2 \left[ 1 - (\hat{\Delta t})^2 \right] \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} \\
 &\quad + (\hat{\Delta t})^2 \left[ \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{jk\Delta z} + \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-jk\Delta z} \right] e^{-j\omega_0 n_z \Delta t} \\
 &= 2 \left[ 1 - (\hat{\Delta t})^2 \right] \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} \\
 &\quad + (\hat{\Delta t})^2 \underbrace{\left[ e^{jk\Delta z} + e^{-jk\Delta z} \right]}_{=2 \cos(k\Delta z)} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} \\
 &= 2 \left[ 1 - (\hat{\Delta t})^2 \right] \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} \\
 &\quad + (\hat{\Delta t})^2 2 \cos(k\Delta z) \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} \\
 \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z+1)\Delta t} &= 2 \left\{ 1 + (\hat{\Delta t})^2 [\cos(k\Delta z) - 1] \right\} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t}
 \end{aligned}$$

$$2 \sin^2 \left( \frac{\alpha}{2} \right) = 1 - \cos \alpha \quad \rightarrow \quad -2 \sin^2 \left( \frac{\alpha}{2} \right) = \cos \alpha - 1$$

$$\begin{aligned}
 \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z+1)\Delta t} &= 2 \left\{ 1 - 2(\hat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} \\
 &= -\hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} + 2 \left\{ 1 - 2(\hat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t}
 \end{aligned}$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

$$\begin{aligned}\hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z+1)\Delta t} &= 2 \left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t} - \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} \\ &= -\hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_z-1)\Delta t} + 2 \left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\} \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_z \Delta t}\end{aligned}$$

Define / Definiere

$$\begin{aligned}U^{(n_t)} &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0 n_t \Delta t} \\ V^{(n_t)} &= U^{(n_t-1)} \\ &= \hat{E}_{0x}(\omega_0, \hat{\mathbf{k}}) e^{-j\omega_0(n_t-1)\Delta t}\end{aligned}$$

which yields for the above equation / womit wir für die obere Gleichung erhalten

$$\begin{aligned}U^{(n_t+1)} &= -U^{(n_t-1)} + 2 \left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\} U^{(n_t)} \\ &= -V^{(n_t)} + 2 \left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\} U^{(n_t)}\end{aligned}$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

Define a new algebraic vector / Definiere einen neuen algebraischen Vektor

$$\{\mathbf{W}\}^{(n_t)} = \begin{Bmatrix} U^{(n_t)} \\ V^{(n_t)} \end{Bmatrix}$$

$$\underbrace{\begin{Bmatrix} U^{(n_t+1)} \\ V^{(n_t+1)} \end{Bmatrix}}_{=\{\mathbf{W}\}^{(n_t+1)}} = \underbrace{\begin{bmatrix} 2\left\{1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right)\right\} & -1 \\ 1 & 0 \end{bmatrix}}_{=[\mathbf{G}]_{1D}^{FD}} \underbrace{\begin{Bmatrix} U^{(n_t)} \\ V^{(n_t)} \end{Bmatrix}}_{=\{\mathbf{W}\}^{(n_t)}}$$

$$\{\mathbf{W}\}^{(n_t+1)} = [\mathbf{G}]_{1D}^{FD} \{\mathbf{W}\}^{(n_t)}$$

$[\mathbf{G}]_{1D}^{FD}$  : Amplification matrix /  
Verstärkungsmatrix

$$\det\left\{[\mathbf{G}]_{1D}^{FD} - \nu[\mathbf{I}]\right\} = 0$$

$\nu_n([\mathbf{G}]_{1D}^{FD})$  :  $n$ th eigenvalue of the matrix  $[\mathbf{G}]_{1D}^{FD}$   
 $n$ -ter Eigenwert der Matrix  $[\mathbf{G}]_{1D}^{FD}$

$$\det\left\{[\mathbf{G}]_{1D}^{FD} - \nu[\mathbf{I}]\right\} = \begin{vmatrix} 2\left\{1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right)\right\} - \nu & -1 \\ 1 & -\nu \end{vmatrix}$$

$$= \nu^2 - 2\nu\left\{1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right)\right\} + 1$$

Characteristic polynomial /  
Charakteristisches Polynom

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

$$v^2 - 2v \left[ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right] + \left[ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right]^2 = \left[ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right]^2 - 1$$

$$\left\{ v - \left[ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right] \right\}^2 = \left[ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right]^2 - 1$$

Eigenvalues of the amplification matrix /  
Eigenwerte der Verstärkungsmatrix

$$v_{1/2} = \underbrace{\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\}}_{=a} \pm \sqrt{\underbrace{\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2 \left( \frac{k\Delta z}{2} \right) \right\}^2}_{=a^2} - 1}$$

$$= a \pm \sqrt{a^2 - 1}$$

**if  $a^2 \leq 1$  / falls  $a^2 \leq 1$**

$$= a \pm j\sqrt{1 - a^2}$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

$$v_{1/2} = \underbrace{\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\}}_{=a} \pm \sqrt{\underbrace{\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\}^2}_{=a^2} - 1}$$

$$= a \pm \sqrt{a^2 - 1}$$

$$= a \pm j\sqrt{1 - a^2} \quad \text{if } a^2 \leq 1 / \text{falls } a^2 \leq 1$$

$$v_n = \text{Re}\{v_n\} + j \text{Im}\{v_n\} \quad n = 1, 2$$

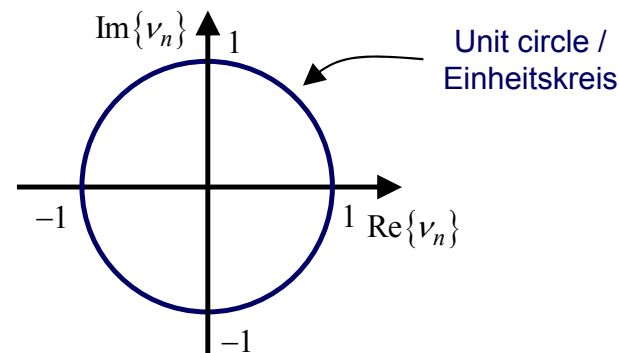
$$|v_{1/2}| = \left| a \pm j\sqrt{1 - a^2} \right| = a^2 + \left( \sqrt{1 - a^2} \right)^2 = a^2 + 1 - a^2 = 1$$



Spectral radius /  
Spektraler Radius

$$\rho\left([\mathbf{G}]_{1D}^{\text{FD}}\right) \leq 1$$

This means for, that all eigenvalues  
 $a^2 \leq 1$  are on the unit circle in the  
complex plane. /  
Dies bedeutet, dass alle Eigenwerte für  
 $a^2 \leq 1$  auf dem Einheitskreis in der  
komplexen Ebene liegen.



## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

$$v_{1/2} = \underbrace{\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\}}_{=a} \pm j \sqrt{1 - \underbrace{\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\}^2}_{=a^2}}$$

$$= a \pm j\sqrt{1-a^2} \quad \text{if } a^2 \leq 1 / \text{ falls } a^2 \leq 1$$

$$a^2 \leq 1$$

$$\left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\}^2 \leq 1$$

$$1 - 4(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) + 4(\widehat{\Delta t})^4 \sin^4\left(\frac{k\Delta z}{2}\right) \leq 1$$

$$-4(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \left[ 1 - (\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right] \leq 0$$

$$1 - (\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \leq 0$$

$$(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \leq 1$$

$$(\widehat{\Delta t})^2 \leq 1 \quad \text{because / weil } \max \left\{ \sin^2\left(\frac{k\Delta z}{2}\right) \right\} = 1$$

$$\widehat{\Delta t} \leq 1$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

1-D Stability Condition for an FD algorithm of 2nd order in space and time– CFL-Condition /  
1D-Stabilitätsbedingung für einen FD-Algorithmus zweiter Ordnung in Raum und Zeit– CFL-Bedingung

$$1\text{-D} / 1\text{D}: \Delta t \leq \Delta t_{\max} = \frac{\Delta x}{c} \quad \widehat{\Delta t} \leq 1$$

2-D and 3-D Stability Condition for an FD algorithm of 2nd order in space and time– CFL-Condition /  
2D- und 3D- Stabilitätsbedingung für einen FD-Algorithmus zweiter Ordnung in Raum und Zeit– CFL-  
Bedingung

$$2\text{-D} / 2\text{D}: \Delta t \leq \Delta t_{\max} = \frac{1}{\sqrt{2}} \frac{\Delta x}{c} \quad \widehat{\Delta t} \leq \frac{1}{\sqrt{2}} \approx 0.707$$

$$3\text{-D} / 3\text{D}: \Delta t \leq \Delta t_{\max} = \frac{1}{\sqrt{3}} \frac{\Delta x}{c} \quad \widehat{\Delta t} \leq \frac{1}{\sqrt{3}} \approx 0.577$$

$$\widehat{\Delta t} = \frac{\Delta t}{\Delta t_{\text{ref}}} : \quad \text{Courant number / Courant - Zahl} \quad \Delta t_{\text{ref}} = \frac{\Delta x}{c}$$

## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

$$v_{1/2}(\widehat{\Delta t}) = a \pm \sqrt{a^2 - 1} \quad \text{if } a^2 \geq 1 / \text{falls } a^2 \geq 1$$

$$= a \pm j\sqrt{1 - a^2} \quad \text{if } a^2 \leq 1 / \text{falls } a^2 \leq 1 \quad \text{with } a = \left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\}$$

$$a^2 \leq 1: v_{1/2} = a \pm j\sqrt{1 - a^2}$$

$$v_1 = a + j\sqrt{1 - a^2} \quad v_2 = a - j\sqrt{1 - a^2}$$

$$|v_1| = 1 \quad |v_2| = 1$$

Spectral radius /  
Spektraler Radius  $\rho([\mathbf{G}]_{1D}^{\text{FD}}) \leq 1$

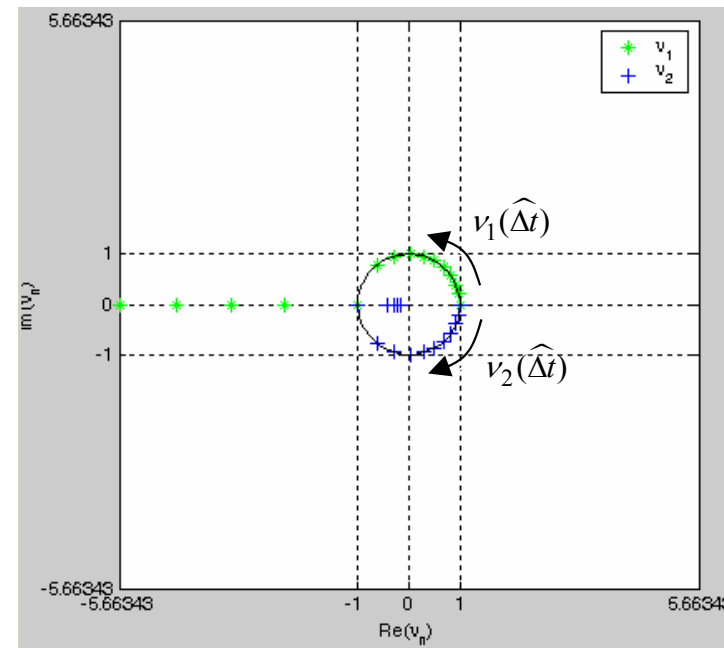
$$a^2 > 1: v_{1/2} = a \pm \sqrt{a^2 - 1}$$

$$v_1 = a + \sqrt{a^2 - 1} \quad v_2 = a - \sqrt{a^2 - 1}$$

$$\lim_{a \rightarrow \infty} |v_1| \rightarrow \infty \quad \lim_{a \rightarrow \infty} |v_2| \rightarrow 0$$

Spectral radius /  
Spektraler Radius  $\rho([\mathbf{G}]_{1D}^{\text{FD}}) \geq 1$

$v_{1/2}(\widehat{\Delta t})$  as a function of  $(\widehat{\Delta t})$



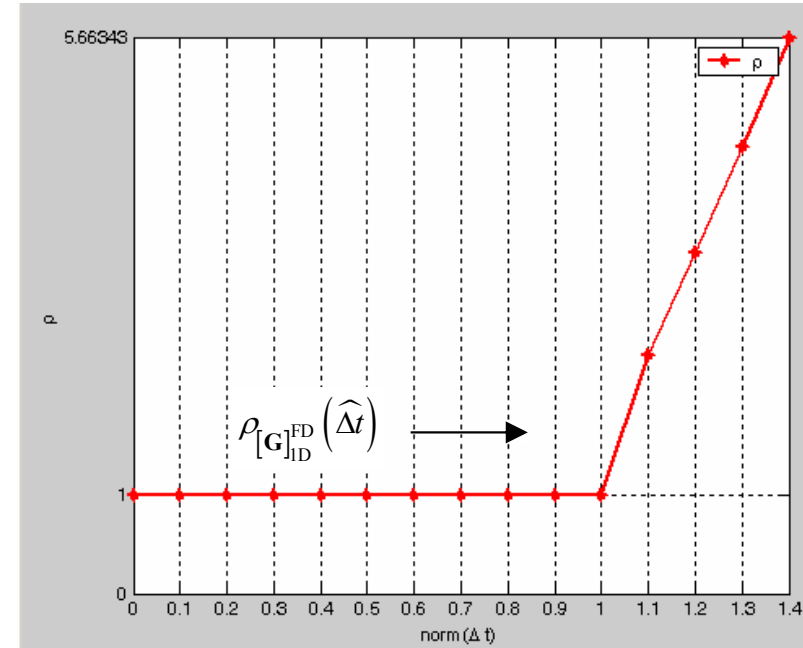
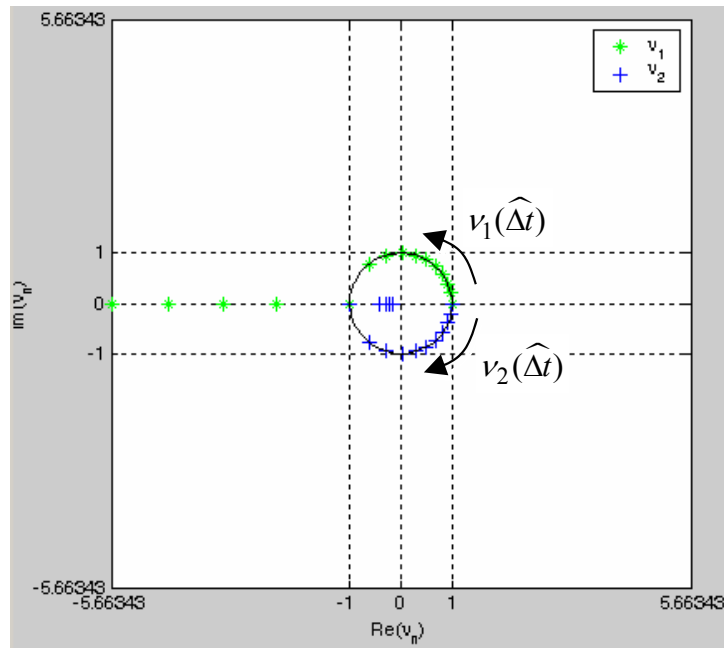
## Derivation of the Stability Condition for the 1-D FD Scheme of 2nd Order / Ableitung der Stabilitätsbedingung für das 1D-FD-Schema 2ter Ordnung

Eigenvalues / Eigenwerte

$$\begin{aligned} v_{1/2}(\widehat{\Delta t}) &= a \pm \sqrt{a^2 - 1} \quad \text{if } a^2 \geq 1 / \text{falls } a^2 \geq 1 \\ &= a \pm j\sqrt{1 - a^2} \quad \text{if } a^2 \leq 1 / \text{falls } a^2 \leq 1 \\ \text{with } a &= \left\{ 1 - 2(\widehat{\Delta t})^2 \sin^2\left(\frac{k\Delta z}{2}\right) \right\} \end{aligned}$$

Spectral radius / Spektraler Radius

$$\rho_{[G]_{1D}}^{\text{FD}}(\widehat{\Delta t})$$



## 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

Maxwell's equations / Maxwell'sche Gleichungen

$$\frac{\partial}{\partial t} \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \rho_e(\underline{\mathbf{R}}, t)$$

Continuity equations / Kontinuitätsgleichungen

$$\nabla \cdot \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_e(\underline{\mathbf{R}}, t)$$

Constitutive Equations for Vacuum /  
Konstituierende Gleichungen  
(Materialgleichungen) für Vakuum

$$\underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \frac{1}{\varepsilon_0} \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\varepsilon_0} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

### 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\frac{\partial}{\partial t} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) \quad (1)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \frac{1}{\varepsilon_0} \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\varepsilon_0} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) \quad (2)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \frac{\partial}{\partial t} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) \quad (3)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \frac{1}{\varepsilon_0} \nabla \times \frac{\partial}{\partial t} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\varepsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) \quad (4)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \left[ \frac{1}{\varepsilon_0} \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\varepsilon_0} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) \right] - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) \quad (5)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \frac{1}{\varepsilon_0} \nabla \times \left[ -\frac{1}{\mu_0} \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) \right] - \frac{1}{\varepsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) \quad (6)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\varepsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) + \frac{1}{\varepsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) \quad (7)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\varepsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\varepsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) - \frac{1}{\varepsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) \quad (8)$$

### 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\frac{1}{\varepsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{H}}(\mathbf{R}, t) + \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\mathbf{R}, t) = + \frac{1}{\varepsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_e(\mathbf{R}, t) - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\mathbf{R}, t) \quad (1)$$

$$\frac{1}{\varepsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{E}}(\mathbf{R}, t) + \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\mathbf{R}, t) = - \frac{1}{\varepsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_m(\mathbf{R}, t) - \frac{1}{\varepsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\mathbf{R}, t) \quad (2)$$

$$c_0 = \frac{1}{\sqrt{\varepsilon_0 \mu_0}}$$

$$-\nabla \times \nabla \times \underline{\mathbf{H}}(\mathbf{R}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\mathbf{R}, t) = -\nabla \times \underline{\mathbf{J}}_e(\mathbf{R}, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\mathbf{R}, t) \quad (3)$$

$$-\nabla \times \nabla \times \underline{\mathbf{E}}(\mathbf{R}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\mathbf{R}, t) = \nabla \times \underline{\mathbf{J}}_m(\mathbf{R}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\mathbf{R}, t) \quad (4)$$

Vector identity /  
Vektoridentität

$$\nabla \times \nabla \times = \nabla \nabla \cdot - \underbrace{\nabla \cdot \nabla}_{=\nabla^2 = \Delta} = \nabla \nabla \cdot - \Delta$$

Short-hand notation /  
Abkürzende Schreibweise

$$\nabla \cdot \nabla = \nabla^2 = \Delta$$

$$-[\nabla \nabla \cdot - \Delta] \underline{\mathbf{H}}(\mathbf{R}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\mathbf{R}, t) = -\nabla \times \underline{\mathbf{J}}_e(\mathbf{R}, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\mathbf{R}, t) \quad (5)$$

$$-[\nabla \nabla \cdot - \Delta] \underline{\mathbf{E}}(\mathbf{R}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\mathbf{R}, t) = \nabla \times \underline{\mathbf{J}}_m(\mathbf{R}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\mathbf{R}, t) \quad (6)$$

## 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\begin{aligned}
 -[\nabla\nabla \cdot - \Delta]\underline{\mathbf{H}}(\underline{\mathbf{R}},t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}},t) &= -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}},t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}},t) \\
 -[\nabla\nabla \cdot - \Delta]\underline{\mathbf{E}}(\underline{\mathbf{R}},t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}},t) &= \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}},t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}},t)
 \end{aligned}$$

3rd and 4th Maxwell's  
equations / 3. und 4.  
Maxwellsche Gleichung

$$\nabla \cdot \underline{\mathbf{B}}(\underline{\mathbf{R}},t) = \rho_m(\underline{\mathbf{R}},t)$$

$$\nabla \cdot \underline{\mathbf{D}}(\underline{\mathbf{R}},t) = \rho_e(\underline{\mathbf{R}},t)$$

Constitutive equations /  
Materialgleichungen

$$\underline{\mathbf{B}}(\underline{\mathbf{R}},t) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}},t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}},t) = \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}},t)$$



$$\nabla \cdot \underline{\mathbf{H}}(\underline{\mathbf{R}},t) = \frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}},t)$$

$$\nabla \cdot \underline{\mathbf{E}}(\underline{\mathbf{R}},t) = \frac{1}{\varepsilon_0} \rho_e(\underline{\mathbf{R}},t)$$

$$\begin{aligned}
 \Delta \underline{\mathbf{H}}(\underline{\mathbf{R}},t) - \nabla \underbrace{\nabla \cdot \underline{\mathbf{H}}(\underline{\mathbf{R}},t)}_{=\frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}},t)} - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}},t) &= -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}},t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}},t) \\
 &= \frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}},t)
 \end{aligned}$$

$$\begin{aligned}
 \Delta \underline{\mathbf{E}}(\underline{\mathbf{R}},t) - \nabla \underbrace{\nabla \cdot \underline{\mathbf{E}}(\underline{\mathbf{R}},t)}_{=\frac{1}{\varepsilon_0} \rho_e(\underline{\mathbf{R}},t)} - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}},t) &= \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}},t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}},t) \\
 &= \frac{1}{\varepsilon_0} \rho_e(\underline{\mathbf{R}},t)
 \end{aligned}$$

### 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \nabla \left[ \frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}}, t) \right] - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \nabla \left[ \frac{1}{\varepsilon_0} \rho_e(\underline{\mathbf{R}}, t) \right] - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \frac{1}{\mu_0} \nabla \rho_m(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \frac{1}{\varepsilon_0} \nabla \rho_e(\underline{\mathbf{R}}, t)$$

Laplace operator in Cartesian coordinates /  
Laplace-Operator in Kartesischen Koordinaten

$$\Delta = \nabla \cdot \nabla$$

$$= \left( \underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot \left( \underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right)$$

$$= \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

## 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\begin{aligned}\Delta \underline{\mathbf{E}}(\mathbf{R}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\mathbf{R}, t)] \\ &= \left( \underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot \left[ \left( \underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \underline{\mathbf{E}}(\mathbf{R}, t) \right]\end{aligned}$$

Short-hand notation /  
Abkürzende Schreibweise  $\frac{\partial}{\partial x} = \partial_x \quad \frac{\partial}{\partial y} = \partial_y \quad \frac{\partial}{\partial z} = \partial_z \quad \frac{\partial}{\partial t} = \partial_t$

$$\begin{aligned}\Delta \underline{\mathbf{E}}(\mathbf{R}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\mathbf{R}, t)] \\ &= (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \cdot [(\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \underline{\mathbf{E}}(\mathbf{R}, t)]\end{aligned}$$

$$\begin{aligned}& (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) [E_x(\mathbf{R}, t) \underline{\mathbf{e}}_x + E_y(\mathbf{R}, t) \underline{\mathbf{e}}_y + E_z(\mathbf{R}, t) \underline{\mathbf{e}}_z] \\ &= \underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\mathbf{R}, t) \\ & \quad + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_y E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\mathbf{R}, t) \\ & \quad + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_z E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_z E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\mathbf{R}, t)\end{aligned}$$

### 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\begin{aligned}
 \Delta \underline{\mathbf{E}}(\mathbf{R}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\mathbf{R}, t)] \\
 &= (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \cdot [(\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \underline{\mathbf{E}}(\mathbf{R}, t)] \\
 &= (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \cdot \left[ \begin{aligned} &\underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_x E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_x E_z(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_y E_z(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\mathbf{R}, t) \end{aligned} \right]
 \end{aligned}$$

$$\begin{aligned}
 \Delta \underline{\mathbf{E}}(\mathbf{R}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\mathbf{R}, t)] \\
 &= \underline{\mathbf{e}}_x \partial_x \cdot \left[ \begin{aligned} &\underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_y E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_z E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_z E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\mathbf{R}, t) \end{aligned} \right] \\
 &+ \underline{\mathbf{e}}_y \partial_y \cdot \left[ \begin{aligned} &\underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_y E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_z E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_z E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\mathbf{R}, t) \end{aligned} \right] \\
 &+ \underline{\mathbf{e}}_z \partial_z \cdot \left[ \begin{aligned} &\underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_y E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\mathbf{R}, t) \\ &+ \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_z E_x(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_z E_y(\mathbf{R}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\mathbf{R}, t) \end{aligned} \right]
 \end{aligned}$$

## 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\begin{aligned}\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)] \\ &= \left[ \underline{\mathbf{e}}_x \partial_x^2 E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \partial_x^2 E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \partial_x^2 E_z(\underline{\mathbf{R}}, t) \right. \\ &\quad + \left[ \underline{\mathbf{e}}_x \partial_y^2 E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \partial_y^2 E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \partial_y^2 E_z(\underline{\mathbf{R}}, t) \right. \\ &\quad \left. \left. + \left[ \underline{\mathbf{e}}_x \partial_z^2 E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \partial_z^2 E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \partial_z^2 E_z(\underline{\mathbf{R}}, t) \right] \right] \\ &= \left( \partial_x^2 + \partial_y^2 + \partial_z^2 \right) \underbrace{\left[ E_x(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_x + E_y(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_y + E_z(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_z \right]}_{=\underline{\mathbf{E}}(\underline{\mathbf{R}}, t)} \\ &= \left( \partial_x^2 + \partial_y^2 + \partial_z^2 \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)\end{aligned}$$

$$\begin{aligned}\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \left( \partial_x^2 + \partial_y^2 + \partial_z^2 \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) \\ &= \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)\end{aligned}$$

### 3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \frac{1}{\mu_0} \nabla \rho_m(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \frac{1}{\varepsilon_0} \nabla \rho_e(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \left( \partial_x^2 + \partial_y^2 + \partial_z^2 \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = \left( \partial_x^2 + \partial_y^2 + \partial_z^2 \right) \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \frac{1}{\mu_0} \nabla \rho_m(\underline{\mathbf{R}}, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \frac{1}{\varepsilon_0} \nabla \rho_e(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_e(\underline{\mathbf{R}}, t)$$

## 2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

We consider the  $xz$  plane and assume that the field is independent of  $y$  /  
Wir betrachten die  $xz$ -Ebene und nehmen an, dass das Feld unabhängig von  $y$  ist ➔  $\frac{\partial}{\partial y} \equiv 0$

Then it follows for the 3-D wave equations / Es folgt dann für die 3D-Wellengleichungen

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = -\nabla \times \underline{\mathbf{J}}_e(x, z, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(x, z, t) + \frac{1}{\mu_0} \nabla \rho_m(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = \nabla \times \underline{\mathbf{J}}_m(x, z, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(x, z, t) + \frac{1}{\varepsilon_0} \nabla \rho_e(x, z, t)$$

And we confine the current sources to / Und wir beschränken die Stromquellen auf

$$\underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = \underline{\mathbf{J}}_m(x, y, t) = J_{my}(x, z, t) \underline{\mathbf{e}}_y$$

$$\underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = \underline{\mathbf{J}}_e(x, y, t) = J_{ey}(x, z, t) \underline{\mathbf{e}}_y$$

This yields for the above given 3-D wave equation /  
Dies ergibt für die oben gegebenen 3D-Wellengleichungen

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = \nabla \times [J_{ey}(x, z, t) \underline{\mathbf{e}}_y] + \varepsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \underline{\mathbf{e}}_y + \frac{1}{\mu_0} \nabla \rho_m(x, y, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = \nabla \times [J_{my}(x, z, t) \underline{\mathbf{e}}_y] + \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \underline{\mathbf{e}}_y + \frac{1}{\varepsilon_0} \nabla \rho_e(x, y, t)$$

## 2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

Curl and divergence of the current sources /  
Rotation und Divergenz der Stromquellen

$$\underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = J_{my}(x, z, t)\underline{\mathbf{e}}_y$$

$$\underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = J_{ey}(x, z, t)\underline{\mathbf{e}}_y$$

$$\begin{aligned} \nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) &= \begin{vmatrix} \underline{\mathbf{e}}_x & \underline{\mathbf{e}}_y & \underline{\mathbf{e}}_z \\ \partial_x & \partial_y & \partial_z \\ 0 & J_{ez}(x, z, t) & 0 \end{vmatrix} \\ &= \frac{\partial}{\partial x} J_{ez}(x, z, t)\underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{ez}(x, z, t)\underline{\mathbf{e}}_x \\ &= -\frac{\partial}{\partial z} J_{ez}(x, z, t)\underline{\mathbf{e}}_x + \frac{\partial}{\partial x} J_{ez}(x, z, t)\underline{\mathbf{e}}_z \end{aligned}$$

$$\begin{aligned} \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) &= \begin{vmatrix} \underline{\mathbf{e}}_x & \underline{\mathbf{e}}_y & \underline{\mathbf{e}}_z \\ \partial_x & \partial_y & \partial_z \\ 0 & J_{mz}(x, z, t) & 0 \end{vmatrix} \\ &= \frac{\partial}{\partial x} J_{mz}(x, z, t)\underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{mz}(x, z, t)\underline{\mathbf{e}}_x \\ &= -\frac{\partial}{\partial z} J_{mz}(x, z, t)\underline{\mathbf{e}}_x + \frac{\partial}{\partial x} J_{mz}(x, z, t)\underline{\mathbf{e}}_z \end{aligned}$$

$$\nabla \cdot J_{my}(x, z, t)\underline{\mathbf{e}}_y = \left( \underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot J_{my}(x, z, t)\underline{\mathbf{e}}_y = \frac{\partial}{\partial y} J_{my}(x, z, t) = 0$$

$$\nabla \cdot J_{ey}(x, z, t)\underline{\mathbf{e}}_y = \left( \underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot J_{ey}(x, z, t)\underline{\mathbf{e}}_y = \frac{\partial}{\partial y} J_{ey}(x, z, t) = 0$$

The divergence of the current sources is in this special case zero, because the currents are constant in  $y$  direction. / Die Divergenz der Stromquellen ist in diesem speziellen Fall null, da die Ströme in  $y$ -Richtung konstant sind.

## 2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

$$\nabla \cdot \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_e(\underline{\mathbf{R}}, t)$$

$$\begin{aligned} \nabla \cdot \underline{\mathbf{J}}_m(x, z, \omega) &= j\omega \rho_m(x, z, \omega) \\ \nabla \cdot \underline{\mathbf{J}}_e(x, z, \omega) &= j\omega \rho_e(x, z, \omega) \end{aligned} \quad \Rightarrow \quad \begin{aligned} \rho_m(x, z, \omega) &= \frac{1}{j\omega} \nabla \cdot \underline{\mathbf{J}}_m(x, z, \omega) \\ \rho_e(x, z, \omega) &= \frac{1}{j\omega} \nabla \cdot \underline{\mathbf{J}}_e(x, z, \omega) \end{aligned}$$

$$\nabla \cdot J_{my}(x, z, \omega) \underline{\mathbf{e}}_y = 0$$

$$\nabla \cdot J_{ey}(x, z, \omega) \underline{\mathbf{e}}_y = 0$$

$$\rho_m(\underline{\mathbf{R}}, \omega) = 0$$

$$\rho_e(\underline{\mathbf{R}}, \omega) = 0$$

$$\Rightarrow \quad \begin{aligned} \rho_m(\underline{\mathbf{R}}, t) &= 0 \\ \rho_e(\underline{\mathbf{R}}, t) &= 0 \end{aligned}$$

## 2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = -\nabla \times \underline{\mathbf{J}}_e(x, z, t) + \varepsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = \nabla \times \underline{\mathbf{J}}_m(x, z, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(x, z, t)$$

$$\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = \frac{\partial}{\partial x} J_{ez}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{ex}(x, z, t) \underline{\mathbf{e}}_x$$

$$\nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = \frac{\partial}{\partial x} J_{mz}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{mx}(x, z, t) \underline{\mathbf{e}}_x$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = -\frac{\partial}{\partial x} J_{ey}(x, z, t) \underline{\mathbf{e}}_z + \frac{\partial}{\partial z} J_{ey}(x, z, t) \underline{\mathbf{e}}_x + \varepsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \underline{\mathbf{e}}_y$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{my}(x, z, t) \underline{\mathbf{e}}_x + \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \underline{\mathbf{e}}_y$$

## 2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = -\frac{\partial}{\partial x} J_{ey}(x, z, t) \underline{\mathbf{e}}_z + \frac{\partial}{\partial z} J_{ey}(x, z, t) \underline{\mathbf{e}}_x + \varepsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \underline{\mathbf{e}}_y$$

Decoupled equations /  
Entkoppelte Gleichungen



$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_x(x, z, t) = \frac{\partial}{\partial z} J_{ey}(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(x, z, t) = \varepsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_z(x, z, t) = -\frac{\partial}{\partial x} J_{ey}(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = -\frac{\partial}{\partial z} J_{my}(x, z, t) \underline{\mathbf{e}}_x + \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \underline{\mathbf{e}}_y + \frac{\partial}{\partial x} J_{my}(x, z, t) \underline{\mathbf{e}}_z$$

Decoupled equations /  
Entkoppelte Gleichungen



$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(x, z, t) = -\frac{\partial}{\partial z} J_{my}(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_z(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t)$$

## 2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall



Separation in 2-D → TM and TE case /  
Separation in 2D → TM- und TE- Fall

TM: transversal magnetic / transversal magnetisch  
TE: transversal electric / transversal elektrisch

TM<sub>y</sub> case / TM<sub>y</sub>-Fall

$$\begin{aligned} E_y(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \\ H_x(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_x(x, z, t) = \frac{\partial}{\partial z} J_{ey}(x, z, t) \\ H_z(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_z(x, z, t) = -\frac{\partial}{\partial x} J_{ey}(x, z, t) \end{aligned}$$

TE<sub>y</sub> case / TE<sub>y</sub>-Fall

$$\begin{aligned} H_y(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(x, z, t) = \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \\ E_x(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(x, z, t) = -\frac{\partial}{\partial z} J_{my}(x, z, t) \\ E_z(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_z(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t) \end{aligned}$$

## 2-D EM Wave Propagation – 2-D TM Case / 2D EM Wellenausbreitung – 2D-TM-Fall

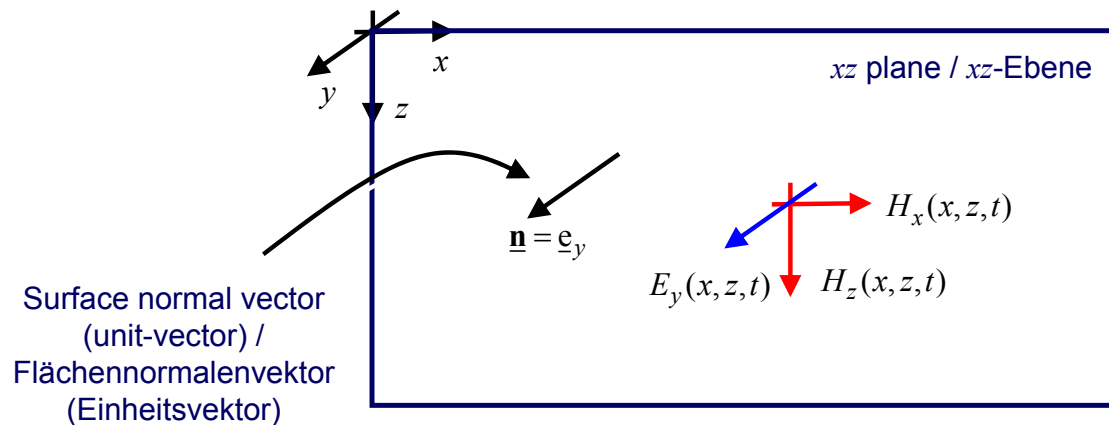


Separation in 2-D → TM case /  
Separation in 2D → TM-Fall

TM: transversal magnetic / transversal magnetisch

TM<sub>y</sub> case / TM<sub>y</sub>-Fall

$$\begin{aligned}
 E_y(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \\
 H_x(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_x(x, z, t) = \frac{\partial}{\partial z} J_{ey}(x, z, t) \\
 H_z(x, z, t) & \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_z(x, z, t) = -\frac{\partial}{\partial x} J_{ey}(x, z, t)
 \end{aligned}$$



## 2-D EM Wave Propagation – 2-D TE Case / 2D EM Wellenausbreitung – 2D-TE-Fall



Separation in 2-D → TE case /  
Separation in 2D → TE- Fall

TE: transversal electric / transversal elektrisch

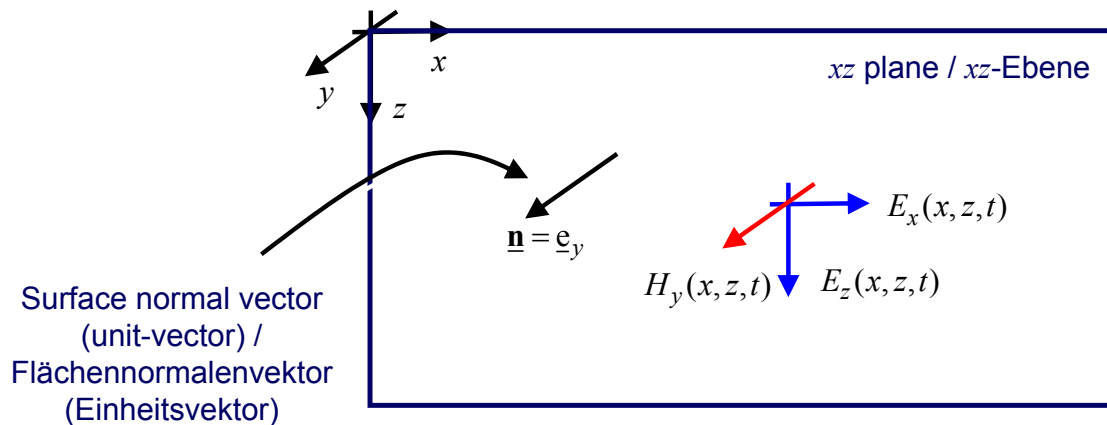
TE<sub>y</sub> case / TE<sub>y</sub>-Fall

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(x, z, t) = \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t)$$

$$H_y(x, z, t) \quad \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(x, z, t) = -\frac{\partial}{\partial z} J_{my}(x, z, t)$$

$$E_x(x, z, t) \quad \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_z(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t)$$

$$E_z(x, z, t)$$



## FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

### Central FD Operators / Zentrale FD-Operatoren

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

### Central FD Operators / Zentrale FD-Operatoren

$$\frac{\partial^2}{\partial x^2} E_y(x, z, t) = \frac{E_y(x + \Delta x, z, t) - 2E_y(x, z, t) + E_y(x - \Delta x, z, t)}{(\Delta x)^2} + O[(\Delta x)^2]$$

$$\frac{\partial^2}{\partial z^2} E_y(x, z, t) = \frac{E_y(x, z + \Delta z, t) - 2E_y(x, z, t) + E_y(x, z - \Delta z, t)}{(\Delta z)^2} + O[(\Delta z)^2]$$

$$\frac{\partial^2}{\partial t^2} E_y(x, z, t) = \frac{E_y(x, z, t + \Delta t) - 2E_y(x, z, t) + E_y(x, z, t - \Delta t)}{(\Delta t)^2} + O[(\Delta t)^2]$$

### Backward FD Operator / Rückwärts-FD-Operator

$$\frac{\partial}{\partial t} J_{ey}(x, z, t) = \frac{J_{ey}(x, z, t) - J_{ey}(x, z, t - \Delta t)}{\Delta t} + O(\Delta t)$$

## FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

### 2-D TM wave equation / 2D-TM-Wellengleichung

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

### Explicit FD algorithm in the time domain of 2nd order in space and time / Expliziter FD-Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit

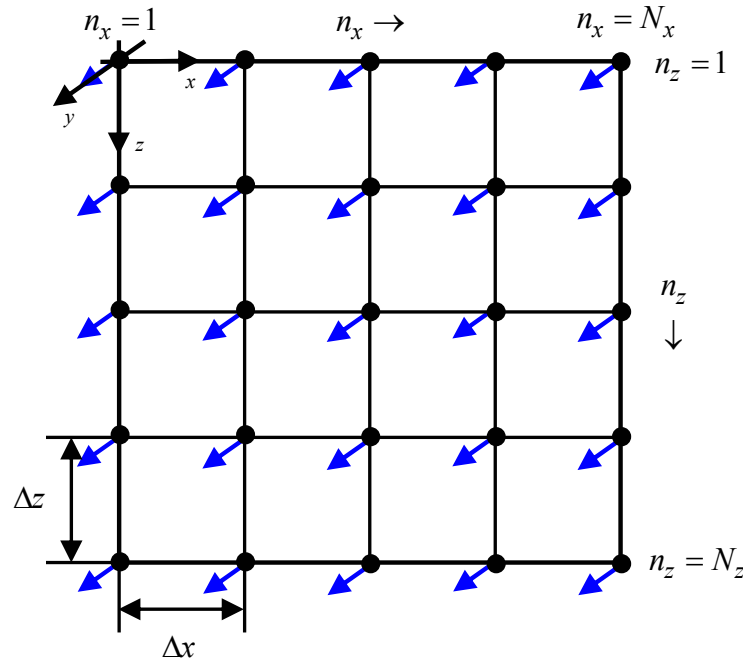
$$\begin{aligned} & \frac{E_y(x + \Delta x, z, t) - 2E_y(x, z, t) + E_y(x - \Delta x, z, t)}{(\Delta x)^2} + \frac{E_y(x, z + \Delta z, t) - 2E_y(x, z, t) + E_y(x, z - \Delta z, t)}{(\Delta z)^2} \\ & - \frac{1}{c_0^2} \frac{E_y(x, z, t + \Delta t) - 2E_y(x, z, t) + E_y(x, z, t - \Delta t)}{(\Delta t)^2} \\ & = \mu_0 \frac{J_{ey}(x, z, t) - J_{ey}(x, z, t - \Delta t)}{\Delta t} + \mathcal{O}[(\Delta x)^2] + \mathcal{O}[(\Delta z)^2] + \mathcal{O}[(\Delta t)^2] \end{aligned}$$

## FD Method – 2-D TM Wave Equation – 2-D FD Grid / FD-Methode – 2D-TM-Wellengleichung – 2D-FD-Gitter

2-D FD grid /  
2D-FD-Gitter

$$E_y(x, z, t) \rightarrow E_y^{(n_x, n_z, n_t)} \rightarrow E_y^{(n, n_t)}$$

■  $E_y(x, z, t) = E_y^{(n_x, n_z, n_t)}$   
 $= E_y^{(n, n_t)}$



$$n_x = 1, \dots, N_x$$

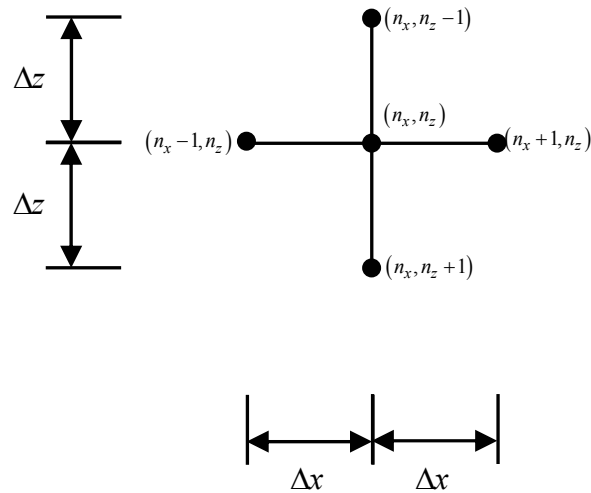
$$n_z = 1, \dots, N_z$$

Global grid node numbering /  
Globale  
Gitterknotennummerierung

$$n = n_x + N_x(n_z - 1) \quad n = 1, \dots, N \quad N = N_x N_z$$

## FD Method – 2-D TM Wave Equation – 2-D FD Stencil / FD-Methode – 2D-TM-Wellengleichung – 2D-FD-Schablone

2-D FD stencil in space /  
2D-FD-Schablone im Raum



## FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

Explicit 2-D FD algorithm in the time domain of 2nd order in space and time /  
Expliziter 2D-FD-Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit

$$\begin{aligned} E_y(x, z, t + \Delta t) &= 2E_y(x, z, t) - E_y(x, z, t - \Delta t) \\ &+ c_0^2 \frac{(\Delta t)^2}{(\Delta x)^2} \left[ E_y(x + \Delta x, z, t) - 2E_y(x, z, t) + E_y(x - \Delta x, z, t) \right] \\ &+ c_0^2 \frac{(\Delta t)^2}{(\Delta z)^2} \left[ E_y(x, z + \Delta z, t) - 2E_y(x, z, t) + E_y(x, z - \Delta z, t) \right] \\ &+ c_0^2 \mu_0 \Delta t \left[ J_{ey}(x, z, t) - J_{ey}(x, z, t - \Delta t) \right] + \mathcal{O}[(\Delta x)^2] + \mathcal{O}[(\Delta z)^2] + \mathcal{O}[(\Delta t)^2] \end{aligned}$$

**Marching-on-in-time algorithm /  
„Marschieren in der Zeit“-Algorithmus**

$$x \rightarrow n_x \Delta x, \quad n_x = 1, \dots, N_x$$

$$z \rightarrow n_z \Delta z, \quad n_z = 1, \dots, N_z$$

$$t \rightarrow n_t \Delta t, \quad n_t = 1, \dots, N_t$$

$$E_y(x, z, t) \rightarrow E_y^{(n_x, n_z, n_t)}$$

$$J_{ey}(x, z, t) \rightarrow J_{ey}^{(n_x, n_z, n_t)}$$

## FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

**Explicit 2-D FD algorithm in the time domain of 2nd order in space and time /  
Expliziter 2D FD-Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit**

$$\begin{aligned}
 E_y^{(n_x, n_z, n_t+1)} &= 2E_y^{(n_x, n_z, n_t)} - E_y^{(n_x, n_z, n_t-1)} \\
 &+ \left( \frac{c_0 \Delta t}{\Delta x} \right)^2 \left[ E_y^{(n_x+1, n_z, n_t)} - 2E_y^{(n_x, n_z, n_t)} + E_y^{(n_x-1, n_z, n_t)} \right] \\
 &+ \left( \frac{c_0 \Delta t}{\Delta z} \right)^2 \left[ E_y^{(n_x, n_z+1, n_t)} - 2E_y^{(n_x, n_z, n_t)} + E_y^{(n_x, n_z-1, n_t)} \right] \\
 &+ c_0^2 \mu_0 \Delta t \left[ J_{ey}^{(n_x, n_z, n_t)} - J_{ey}^{(n_x, n_z, n_t-1)} \right] + \mathcal{O}[(\Delta x)^2] + \mathcal{O}[(\Delta z)^2] + \mathcal{O}[(\Delta t)^2]
 \end{aligned}$$

**Homogeneous 2-D FD grid of quadratic cells /  
Homogenes 2D- FD-Gitter aus quadratischen Zellen**       $\Delta x = \Delta z$

$$\begin{aligned}
 E_y^{(n_x, n_z, n_t+1)} &= 2E_y^{(n_x, n_z, n_t)} - E_y^{(n_x, n_z, n_t-1)} \\
 &+ \left( \frac{c_0 \Delta t}{\Delta x} \right)^2 \left[ E_y^{(n_x, n_z+1, n_t)} + E_y^{(n_x+1, n_z, n_t)} - 4E_y^{(n_x, n_z, n_t)} + E_y^{(n_x-1, n_z, n_t)} + E_y^{(n_x, n_z-1, n_t)} \right] \\
 &+ c_0^2 \mu_0 \Delta t \left[ J_{ey}^{(n_x, n_z, n_t)} - J_{ey}^{(n_x, n_z, n_t-1)} \right] + \mathcal{O}[(\Delta x)^2] + \mathcal{O}[(\Delta t)^2]
 \end{aligned}$$

## FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

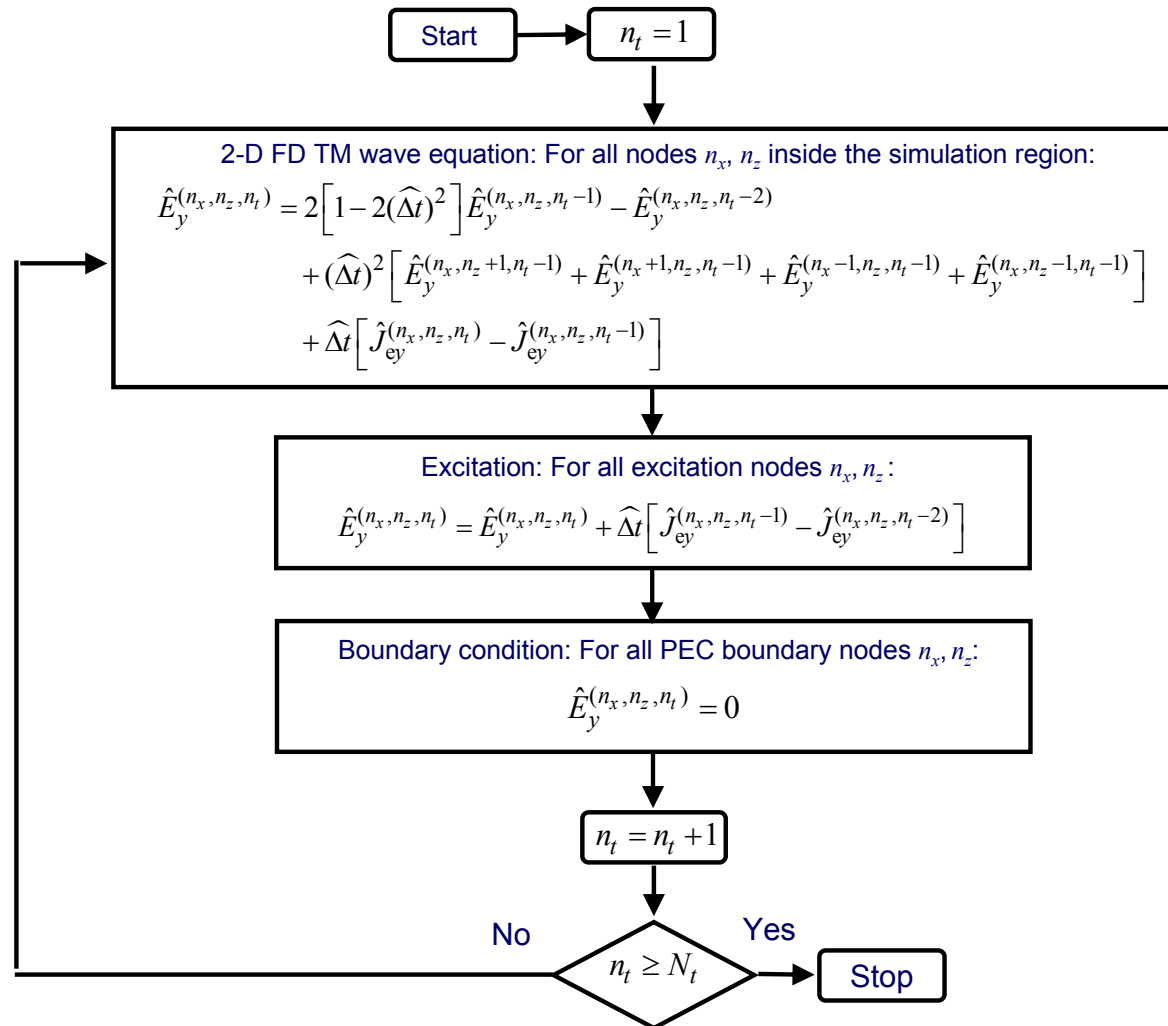
**Explicit FD algorithm in the time domain of 2nd order in space and time /  
Expliziter FD-Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit**

$$\begin{aligned} \hat{E}_y^{(n_x, n_z, n_t+1)} &= 2\hat{E}_y^{(n_x, n_z, n_t)} - \hat{E}_y^{(n_x, n_z, n_t-1)} \\ &+ (\widehat{\Delta t})^2 \left[ \hat{E}_y^{(n_x, n_z+1, n_t)} + \hat{E}_y^{(n_x+1, n_z, n_t)} - 4\hat{E}_y^{(n_x, n_z, n_t)} + \hat{E}_y^{(n_x-1, n_z, n_t)} + \hat{E}_y^{(n_x, n_z-1, n_t)} \right] \\ &+ \widehat{\Delta t} \left[ \hat{J}_{ey}^{(n_x, n_z, n_t)} - \hat{J}_{ey}^{(n_x, n_z, n_t-1)} \right] \end{aligned}$$

$$\text{for / für } \begin{cases} 1 \leq n_x \leq N_x \\ 1 \leq n_z \leq N_z \\ 1 \leq n_t \leq N_t \end{cases}$$

$$\begin{aligned} \hat{E}_y^{(n_x, n_z, n_t)} &= 2 \left[ 1 - 2(\widehat{\Delta t})^2 \right] \hat{E}_y^{(n_x, n_z, n_t-1)} - \hat{E}_y^{(n_x, n_z, n_t-2)} \\ &+ (\widehat{\Delta t})^2 \left[ \hat{E}_y^{(n_x, n_z+1, n_t)} + \hat{E}_y^{(n_x+1, n_z, n_t)} + \hat{E}_y^{(n_x-1, n_z, n_t)} + \hat{E}_y^{(n_x, n_z-1, n_t)} \right] \\ &+ \widehat{\Delta t} \left[ \hat{J}_{ey}^{(n_x, n_z, n_t)} - \hat{J}_{ey}^{(n_x, n_z, n_t-1)} \right] \end{aligned}$$

## FD Method – 2-D FD Wave Equation – TM Case – Flow Chart / FD-Methode – 2D FD-Wellengleichung – TM-Fall – Flussdiagramm



## FD Method – 2-D TM Wave Equation – Example / FD-Methode – 2D-TM-Wellengleichung – Beispiel

Scalar 2-D TM wave equation / Skalare 2D-TM-Wellengleichung

$$\left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \quad \text{for / für} \quad \begin{cases} 0 \leq x \leq X \\ 0 \leq z \leq Z \\ 0 \leq t \leq T \end{cases}$$

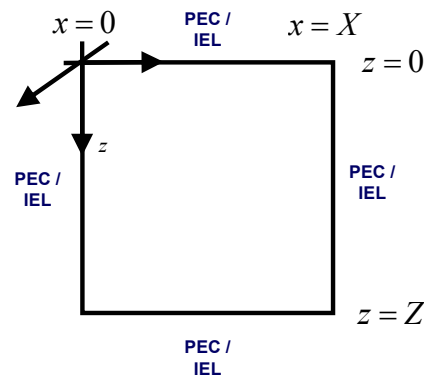
Initial condition / Anfangsbedingung

$$\begin{aligned} E_y(x, z, t) = J_{ey}(x, z, t) = 0 & \quad t \leq 0 \\ J_{ey}(x, z, t) = \delta(x - x_0) \delta(z - z_0) f(t) & \quad t > 0 \end{aligned} \quad \text{Causality / Kausalität}$$

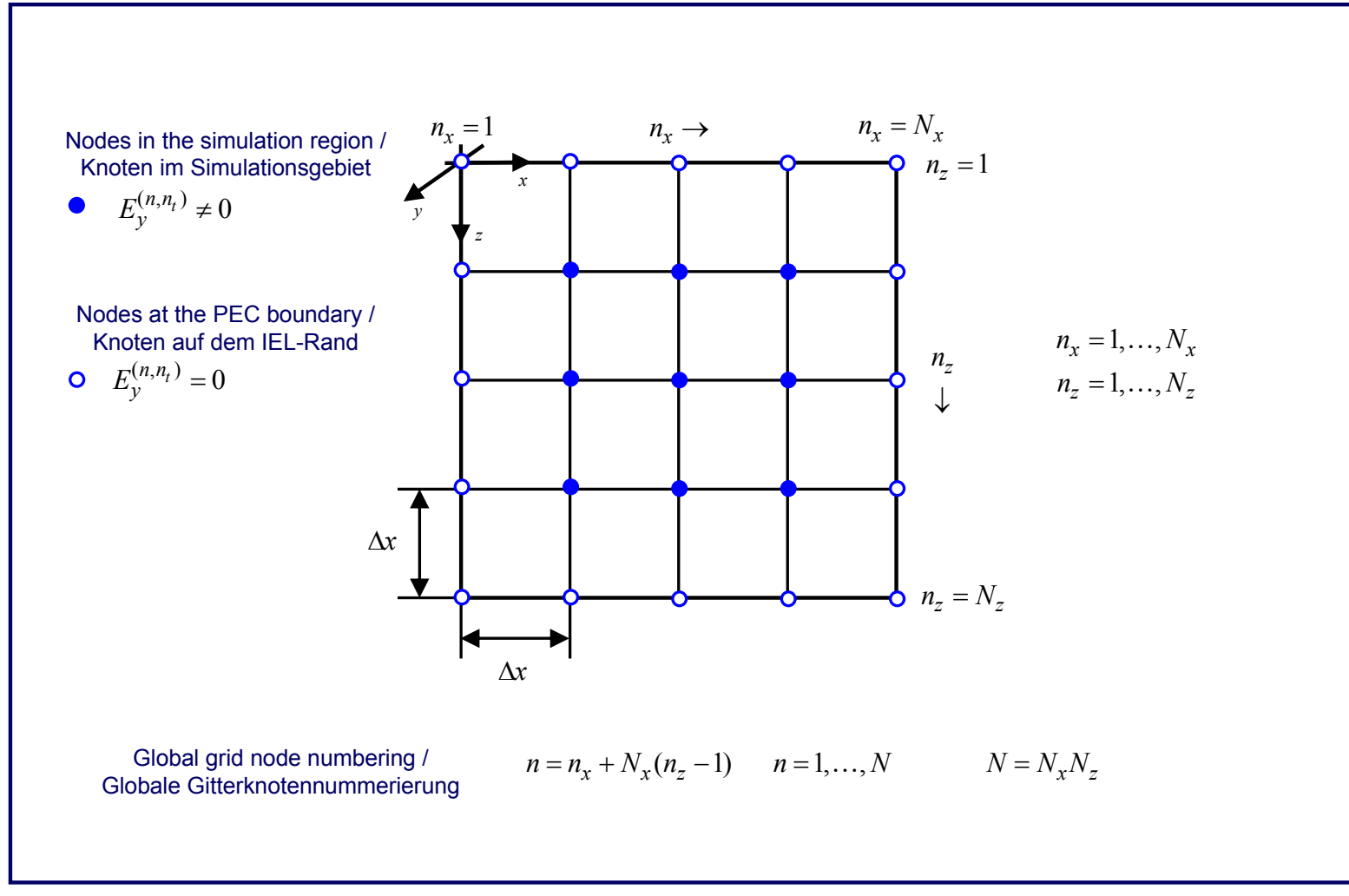
Boundary conditions for a perfectly electrically conducting (PEC) boundary /  
Randbedingung für einen ideal elektrisch leitenden (IEL) Rand

$$\left. \begin{aligned} E_y(0, z, t) = 0 \\ E_y(X, z, t) = 0 \end{aligned} \right\} \forall z, t \forall t \quad \text{and / und} \quad \left. \begin{aligned} E_y(x, 0, t) = 0 \\ E_y(x, Z, t) = 0 \end{aligned} \right\} \forall x, t \forall t$$

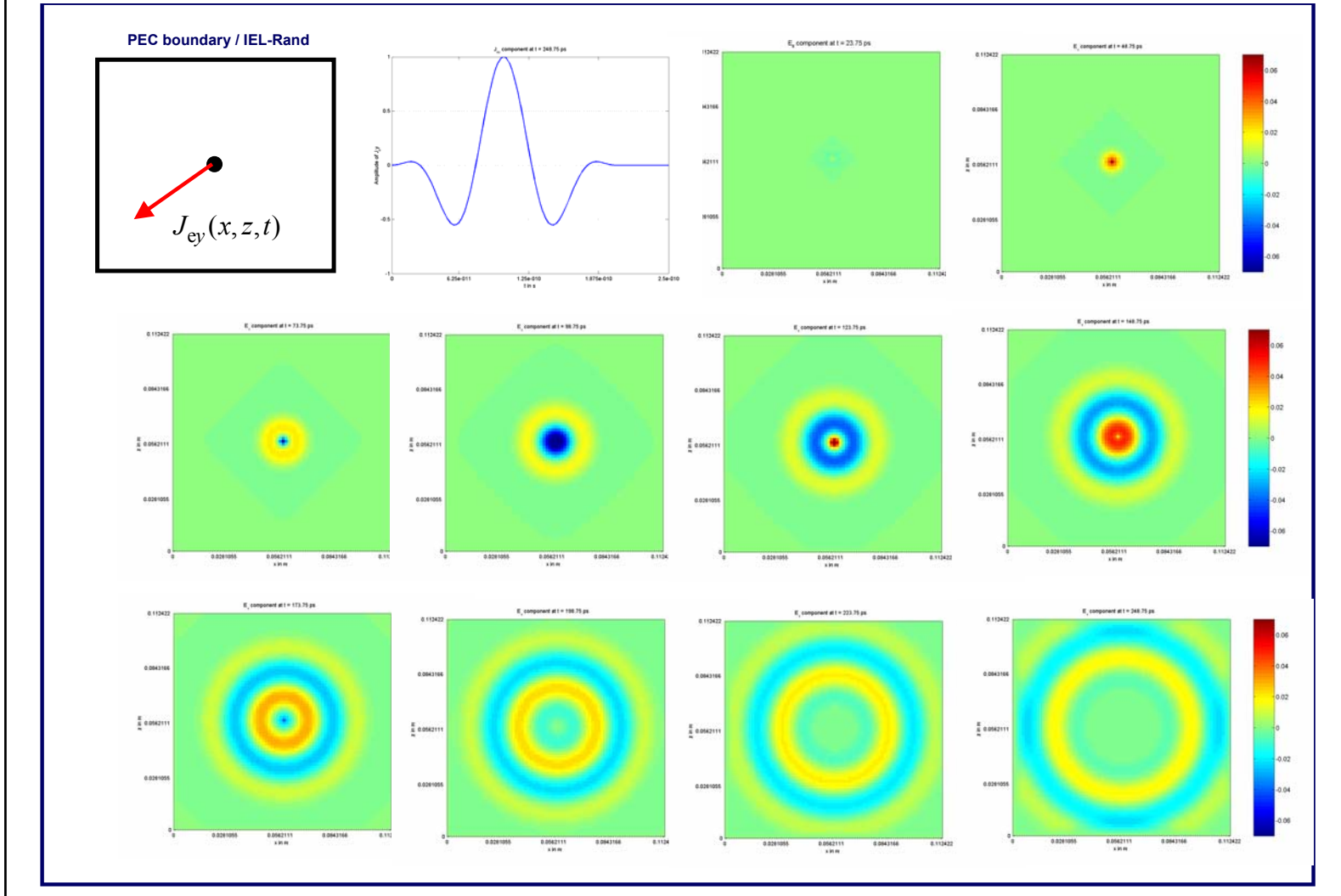
Hyperbolic initial-  
boundary-value  
problem /  
Hyperbolisches  
Anfangs-Randwert-  
Problem



## FD Method – 2-D TM Wave Equation – Example / FD-Methode – 2D-TM-Wellengleichung – Beispiel



# FD Method – 2-D TM Wave Equation – Example/ FD-Methode – 2D-TM-Wellengleichung – Beispiel



## FD Method – 2-D FD Wave Equation – TM Case – Validation / FD-Methode – 2D FD-Wellengleichung – TM-Fall – Validierung

Domain integral representation /  
(Gebiets-) Integraldarstellung

1-D case / 1D-Fall

$$E_x(z, \omega) = j\omega\mu_0 \int_{z'=-\infty}^{\infty} G(z-z', \omega) J_{ex}(z', \omega) dz'$$

2-D case / 2D-Fall

$$E_y(x, z, \omega) = j\omega\mu_0 \int_{z'=-\infty}^{\infty} \int_{x'=-\infty}^{\infty} G(x-x', z-z', \omega) J_{ey}(x', z', \omega) dx' dz'$$

Green's function / Greensche Funktion

$$G(z, \omega) = \frac{c_0}{2} \left[ j \text{PV} \frac{1}{\omega_0} + \pi \delta(z) \right] e^{jk_0|z|}$$

$$G(z, t) = \frac{c_0}{2} u \left( t - \frac{|z|}{c_0} \right)$$

$$G(x, z, \omega) = \frac{j}{4} H_0^{(1)} \left( \frac{\omega}{c_0} \sqrt{x^2 + z^2} \right)$$

$$G(x, z, t) = \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - (x^2 + z^2)}} u \left( t - \frac{\sqrt{x^2 + z^2}}{c_0} \right)$$

## FD Method – 2-D FD Wave Equation – TM Case – Validation / FD-Methode – 2D FD-Wellengleichung – TM-Fall – Validierung

2-D Domain integral representation /  
2D-(Gebiets-)Integraldarstellung

$$E_y(r, \omega) = j\omega\mu_0 \int_{r'=0}^{\infty} G(r-r', \omega) J_{ey}(r', \omega) dr'$$

$$G(r, \omega) = \frac{j}{4} H_0^{(1)} \left( \frac{\omega}{c_0} r \right)$$

$$G(r, t) = \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - r^2}} u \left( t - \frac{r}{c_0} \right)$$

$$E_y(\mathbf{r}, \omega) = j\omega\mu_0 \int_{\mathbf{r}'} G(\mathbf{r}-\mathbf{r}', \omega) J_{ey}(\mathbf{r}', \omega) d\mathbf{r}'$$

$$G(\mathbf{r}-\mathbf{r}', \omega) = \frac{j}{4} H_0^{(1)} \left( \frac{\omega}{c_0} |\mathbf{r}-\mathbf{r}'| \right)$$

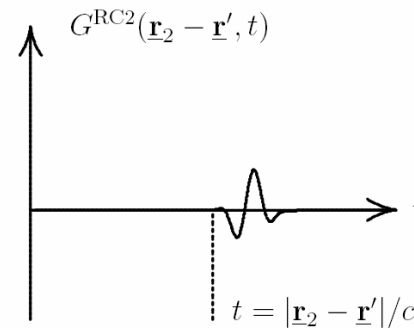
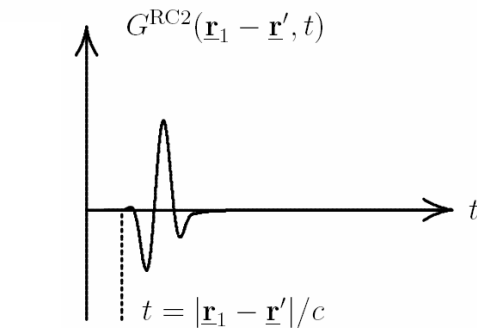
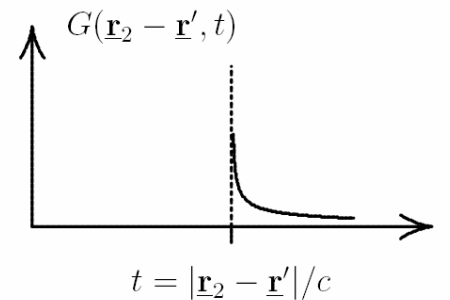
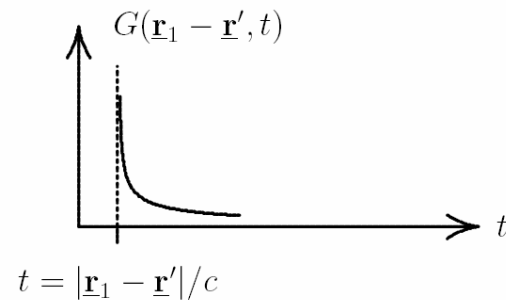
$$G(\mathbf{r}-\mathbf{r}', t) = \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - |\mathbf{r}-\mathbf{r}'|^2}} u \left( t - \frac{|\mathbf{r}-\mathbf{r}'|}{c_0} \right)$$

## FD Method – 2-D FD Wave Equation – TM Case – Validation / FD-Methode – 2D FD-Wellengleichung – TM-Fall – Validierung

2-D Domain integral representation / 2D-(Gebiets-)Integraldarstellung

$$G^{\text{RC2}}(\underline{\mathbf{r}} - \underline{\mathbf{r}}', \omega) = \text{RC2}(\omega) \frac{j}{4} H_0^{(1)} \left( \frac{\omega}{c_0} |\underline{\mathbf{r}} - \underline{\mathbf{r}}'| \right)$$

$$G^{\text{RC2}}(\underline{\mathbf{r}} - \underline{\mathbf{r}}', t) = \text{RC2}(t) * \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - |\underline{\mathbf{r}} - \underline{\mathbf{r}}'|^2}} u \left( t - \frac{|\underline{\mathbf{r}} - \underline{\mathbf{r}}'|}{c_0} \right)$$



$$G^{\text{RC2}}(\underline{\mathbf{r}} - \underline{\mathbf{r}}', \omega) \approx \frac{1}{4} e^{j\frac{\pi}{4}} \sqrt{\frac{2c}{\pi}} \frac{\text{RC2}(\omega)}{\sqrt{\omega}} \frac{e^{jk|\underline{\mathbf{r}} - \underline{\mathbf{r}}'|}}{\sqrt{|\underline{\mathbf{r}} - \underline{\mathbf{r}}'|}}$$

**End of Lecture 4 /  
Ende der 4. Vorlesung**