

Elektromagnetische Feldtheorie I (EFT I) / Electromagnetic Field Theory I (EFT I)

10th Lecture / 10. Vorlesung

Dr.-Ing. René Marklein

marklein@uni-kassel.de

<http://www.tet.e-technik.uni-kassel.de>

<http://www.uni-kassel.de/fb16/tet/marklein/index.html>

Universität Kassel

Fachbereich Elektrotechnik / Informatik

(FB 16)

Fachgebiet Theoretische Elektrotechnik

(FG TET)

Wilhelmshöher Allee 71

Büro: Raum 2113 / 2115

D-34121 Kassel

University of Kassel

**Dept. Electrical Engineering / Computer
Science (FB 16)**

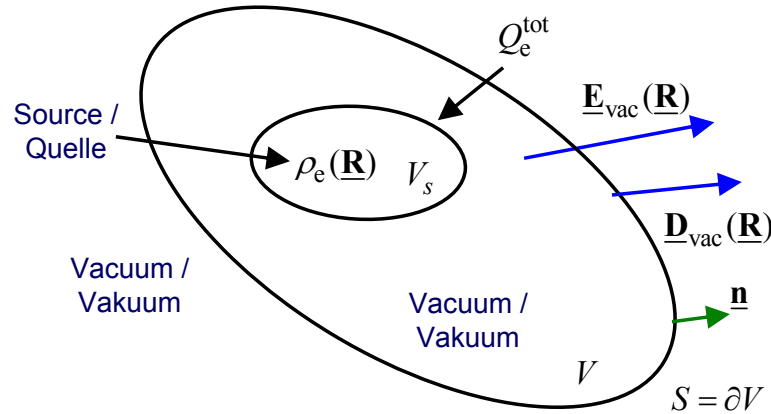
**Electromagnetic Field Theory
(FG TET)**

Wilhelmshöher Allee 71

Office: Room 2113 / 2115

D-34121 Kassel

ES Fields – Electric Polarization of Materials / ES Felder – Elektrische Polarisation von Materialien

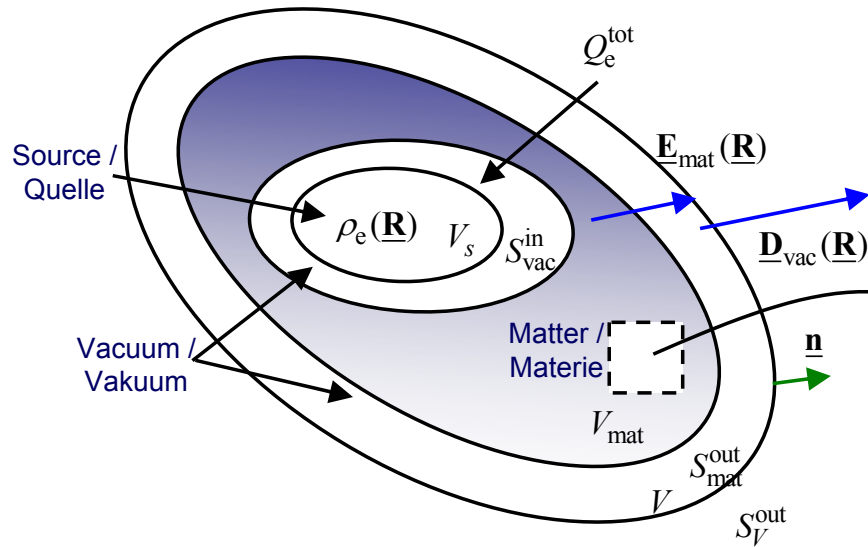


$$\oiint_{S_s = \partial V_s} \mathbf{D}_{\text{vac}}(\mathbf{R}) \cdot \mathbf{n} dS = Q_e^{\text{tot}}$$

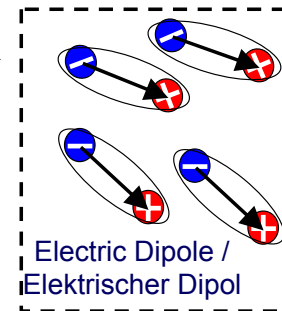
$$= \oiint_{S = \partial V} \mathbf{D}_{\text{vac}}(\mathbf{R}) \cdot \mathbf{n} dS$$

$$\mathbf{D}_{\text{vac}}(\mathbf{R}) = \epsilon_0 \mathbf{E}_{\text{vac}}(\mathbf{R}) \quad \text{Vacuum / Vakuum}$$

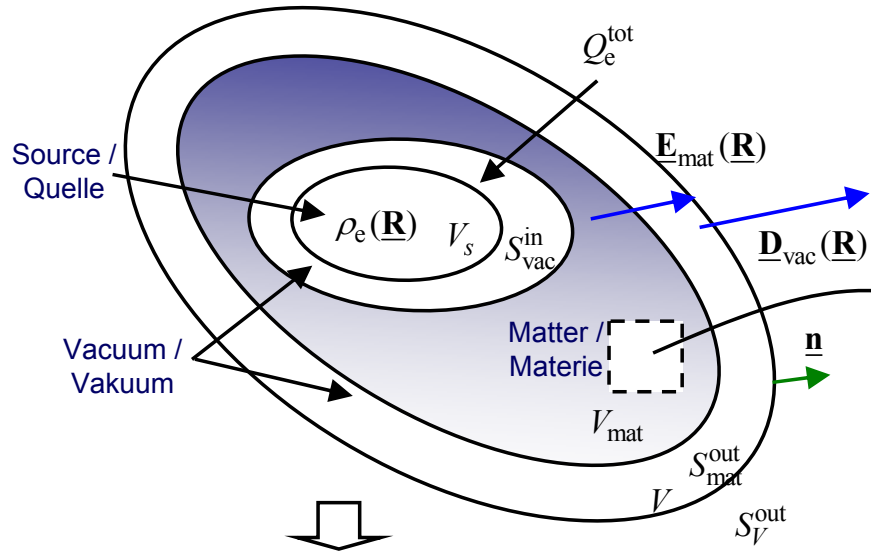
$$= \epsilon_0 \underbrace{\epsilon_{r,\text{vac}}}_{=1} \mathbf{E}_{\text{vac}}(\mathbf{R}) \quad \epsilon_{r,\text{vac}} = 1$$



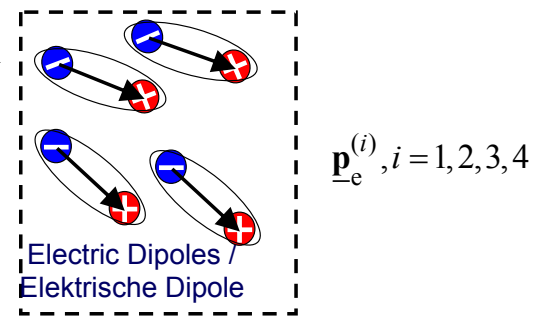
Because of the applied E Field, which is generated by the source, the matter exhibits an electric polarization. /
Aufgrund des anliegenden E-Feldes, welches durch die Quelle generiert wird, kommt es zur elektrischen Polarisation der Materie



ES Fields – Electric Polarization of Materials / ES Felder – Elektrische Polarisierung von Materialien



Because of the applied E Field, which is generated by the source, the matter exhibits an electric polarization. /
Aufgrund des anliegenden E-Feldes, welches durch die Quelle generiert wird, kommt es zur elektrischen Polarisierung der Materie



Total Electric Dipole Density / Elektrisches Gesamtdipolmoment

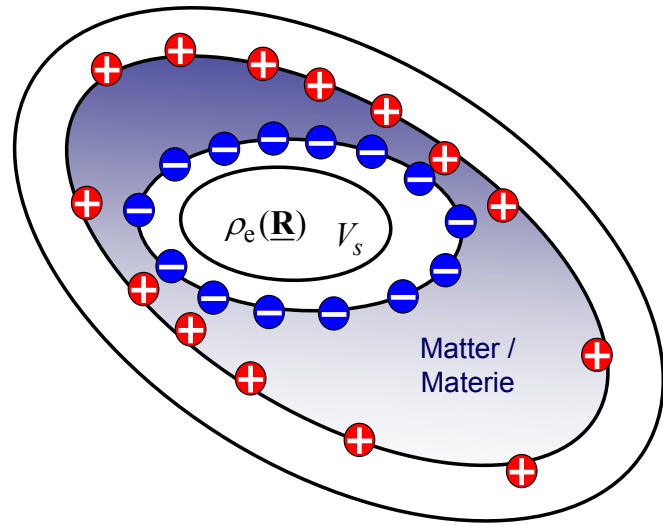
$$\iiint_{V_M} \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) dV = \sum_{i=1}^N \underline{\mathbf{p}}_e^{(i)}$$

Electric Dipole Moment Density / Elektrische Dipolmomentendichte

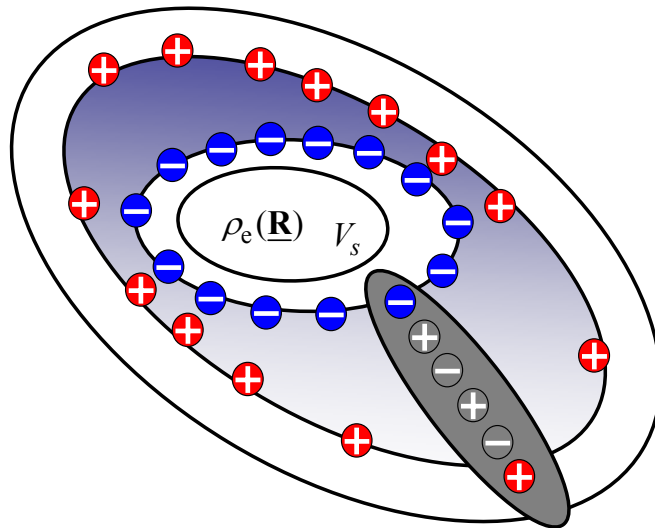
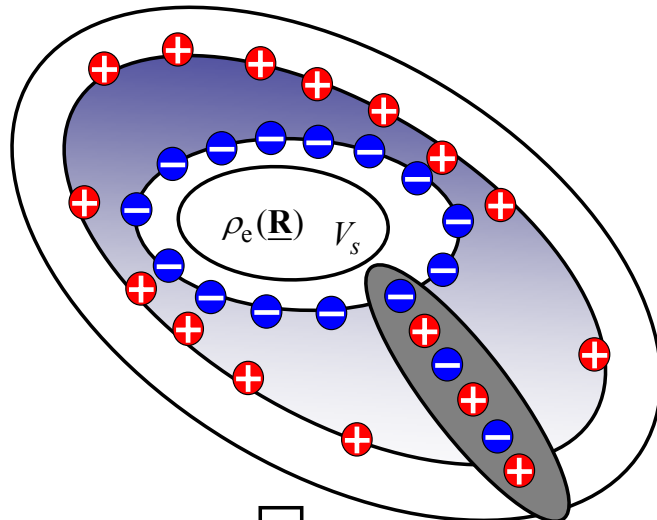
$$\underline{\boldsymbol{\pi}}_e^{(i)}(\underline{\mathbf{R}}) = \underline{\mathbf{p}}_e^{(i)} \delta(\underline{\mathbf{R}} - \underline{\mathbf{R}}^{(i)})$$

Total Electric Dipole Moment Density / Elektrische Gesamtdipolmomentendichte

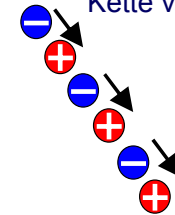
$$\underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = \sum_{i=1}^N \underline{\boldsymbol{\pi}}_e^{(i)}(\underline{\mathbf{R}})$$



ES Fields – Electric Polarization of Materials / ES Felder – Elektrische Polarisierung von Materialien



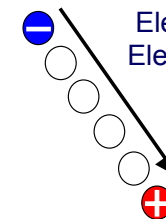
Chain of Electric Dipoles /
Kette von elektrischen Dipolen



Internal Electric Dipoles
Compensate /
Innere elektrische Dipole
kompensieren sich



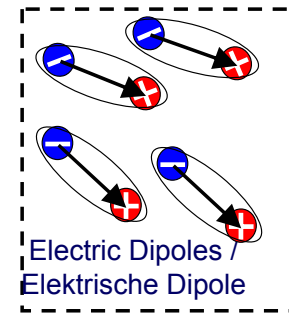
Electric Surface Charge /
Elektrische Flächenladung



ES Fields – Electric Polarization of Materials / ES Felder – Elektrische Polarisation von Materialien

$$\iiint_{V_M} \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) dV = \sum_{i=1}^N \underline{\mathbf{p}}_e^{(i)}$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}) = \begin{cases} \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) & \underline{\mathbf{R}} \in V_{\text{mat}} \\ \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) & \underline{\mathbf{R}} \in V_{\text{vac}} \end{cases}$$



$\underline{\mathbf{p}}_e^{(i)}, i = 1, 2, 3, 4$

$$\begin{aligned} \underline{\mathbf{D}}(\underline{\mathbf{R}}) &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \varepsilon_r(\underline{\mathbf{R}}) \underline{\mathbf{E}}(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underbrace{\varepsilon_0 [\varepsilon_r(\underline{\mathbf{R}}) - 1] \underline{\mathbf{E}}(\underline{\mathbf{R}})}_{=\underline{\mathbf{P}}_e(\underline{\mathbf{R}})} \end{aligned}$$

$\varepsilon_r(\underline{\mathbf{R}})$ Relative Permittivity /
Relative Permittivität

$$\begin{aligned} \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) &= \varepsilon_0 \underbrace{[\varepsilon_r(\underline{\mathbf{R}}) - 1]}_{=\chi_e(\underline{\mathbf{R}})} \underline{\mathbf{E}}(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \chi_e(\underline{\mathbf{R}}) \underline{\mathbf{E}}(\underline{\mathbf{R}}) \end{aligned}$$

$\chi_e(\underline{\mathbf{R}})$ Electric Susceptibility /
Elektrische Suszeptibilität

ES Fields – Electric Polarization of Materials / ES Felder – Elektrische Polarisation von Materialien

General Case / Allgemeiner Fall

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}) = \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underline{\mathbf{P}}_e(\underline{\mathbf{R}})$$

Isotropic Case / Isotroper Fall

$$\underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = \varepsilon_0 \chi_e(\underline{\mathbf{R}}) \underline{\mathbf{E}}(\underline{\mathbf{R}})$$

$$\begin{aligned} \underline{\mathbf{D}}(\underline{\mathbf{R}}) &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \varepsilon_0 \chi_e(\underline{\mathbf{R}}) \underline{\mathbf{E}}(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \varepsilon_0 [\varepsilon_r(\underline{\mathbf{R}}) - 1] \underline{\mathbf{E}}(\underline{\mathbf{R}}) \end{aligned}$$

$$\chi_e(\underline{\mathbf{R}}) = \varepsilon_r(\underline{\mathbf{R}}) - 1$$

$$\varepsilon_r(\underline{\mathbf{R}}) = \chi_e(\underline{\mathbf{R}}) + 1$$

Anisotropic Case / Anisotroper Fall

$$\underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = \varepsilon_0 \underline{\underline{\chi}}_e(\underline{\mathbf{R}}) \cdot \underline{\mathbf{E}}(\underline{\mathbf{R}})$$

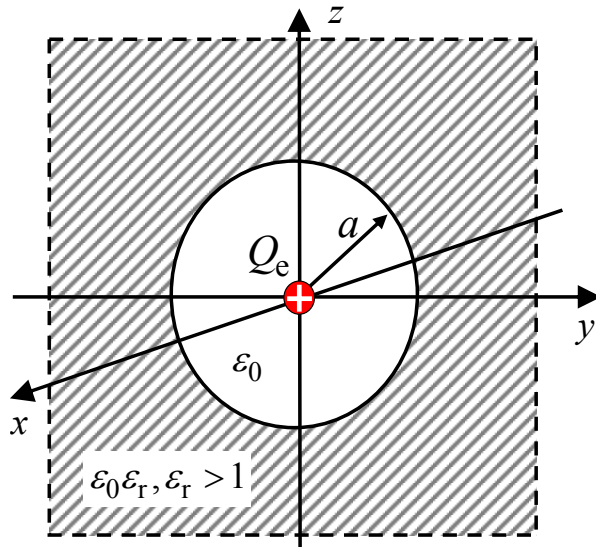
$$\begin{aligned} \underline{\mathbf{D}}(\underline{\mathbf{R}}) &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \varepsilon_0 \underline{\underline{\chi}}_e(\underline{\mathbf{R}}) \cdot \underline{\mathbf{E}}(\underline{\mathbf{R}}) \\ &= \varepsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \varepsilon_0 [\underline{\underline{\varepsilon}}_r(\underline{\mathbf{R}}) - \underline{\underline{\mathbf{I}}}] \cdot \underline{\mathbf{E}}(\underline{\mathbf{R}}) \end{aligned}$$

$$\underline{\underline{\chi}}_e(\underline{\mathbf{R}}) = \underline{\underline{\varepsilon}}_r(\underline{\mathbf{R}}) - \underline{\underline{\mathbf{I}}}$$

$$\underline{\underline{\varepsilon}}_r(\underline{\mathbf{R}}) = \underline{\underline{\chi}}_e(\underline{\mathbf{R}}) + \underline{\underline{\mathbf{I}}}$$

ES Fields – Electric Polarization of Materials – Example / ES Felder – Elektrische Polarisierung von Materialien – Beispiel

Electric Point Charge Embedded in a Sphere Filled with Vacuum, which is Embedded in a Dielectric Material /
Elektrische Punktladung eingebettet in einer mit Vakuum gefüllter Kugel, die in ein dielektrisches Material eingebettet ist.



$$\epsilon(R) = \begin{cases} \epsilon_0 & R < a \\ \epsilon_0 \epsilon_r & R > a \end{cases} \quad \underline{\mathbf{D}}(\underline{\mathbf{R}}) = \frac{Q_e}{4\pi} \frac{1}{R^2} \hat{\underline{\mathbf{R}}} = \frac{Q_e}{4\pi} \frac{1}{R^2} \underline{\mathbf{e}}_R$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \begin{cases} \frac{\underline{\mathbf{D}}(\underline{\mathbf{R}})}{\epsilon_0} & R < a \\ \frac{\underline{\mathbf{D}}(\underline{\mathbf{R}})}{\epsilon_0 \epsilon_r} & R > a \end{cases}$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}) = \begin{cases} \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) & R < a \\ \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}) + \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) & R > a \end{cases}$$

$$\underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = \begin{cases} \underline{\mathbf{0}} & R < a \\ \epsilon_0 (\epsilon_r - 1) \underline{\mathbf{E}}(\underline{\mathbf{R}}) & R > a \end{cases}$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \begin{cases} \frac{\underline{\mathbf{D}}(\underline{\mathbf{R}})}{\epsilon_0} & R < a \\ \frac{\underline{\mathbf{D}}(\underline{\mathbf{R}})}{\epsilon_0} + \frac{\underline{\mathbf{P}}_e(\underline{\mathbf{R}})}{\epsilon_0} & R > a \end{cases} = \begin{cases} \frac{Q_e}{4\pi \epsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R & R < a \\ \frac{Q_e}{4\pi \epsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R - \frac{\underline{\mathbf{P}}_e(\underline{\mathbf{R}})}{\epsilon_0} & R > a \end{cases}$$

ES Fields – Electric Polarization of Materials – Example / ES Felder – Elektrische Polarisierung von Materialien – Beispiel (...)

$$\underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = \begin{cases} \underline{\mathbf{0}} & R < a \\ \varepsilon_0(\varepsilon_r - 1)\underline{\mathbf{E}}(\underline{\mathbf{R}}) & R > a \end{cases} \quad R > a$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R - (\varepsilon_r - 1)\underline{\mathbf{E}}(\underline{\mathbf{R}})$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \begin{cases} \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R & R < a \\ \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R - \frac{\underline{\mathbf{P}}_e(\underline{\mathbf{R}})}{\varepsilon_0} & R > a \end{cases}$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) + (\varepsilon_r - 1)\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R$$

$$\varepsilon_r \underline{\mathbf{E}}(\underline{\mathbf{R}}) = \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R$$

$$= \begin{cases} \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R & R < a \\ \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R - \frac{\varepsilon_0(\varepsilon_r - 1)\underline{\mathbf{E}}(\underline{\mathbf{R}})}{\varepsilon_0} & R > a \end{cases}$$

$$= \begin{cases} \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R & R < a \\ \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R - (\varepsilon_r - 1)\underline{\mathbf{E}}(\underline{\mathbf{R}}) & R > a \end{cases}$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \begin{cases} \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R & R < a \\ \frac{Q_e}{4\pi\varepsilon_0\varepsilon_r} \frac{1}{R^2} \underline{\mathbf{e}}_R & R > a \end{cases}$$

ES Fields – Electric Polarization of Materials – Example / ES Felder – Elektrische Polarisation von Materialien – Beispiel (...)

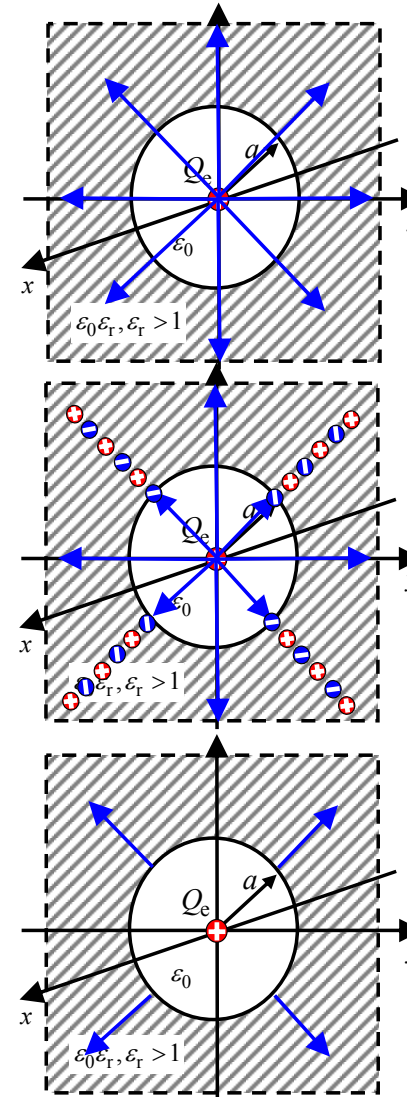
**Permittivity Profile /
Permittivitätsprofil**

$$\varepsilon(R) = \begin{cases} \varepsilon_0 & R < a \\ \varepsilon_0 \varepsilon_r & R > a \end{cases}$$

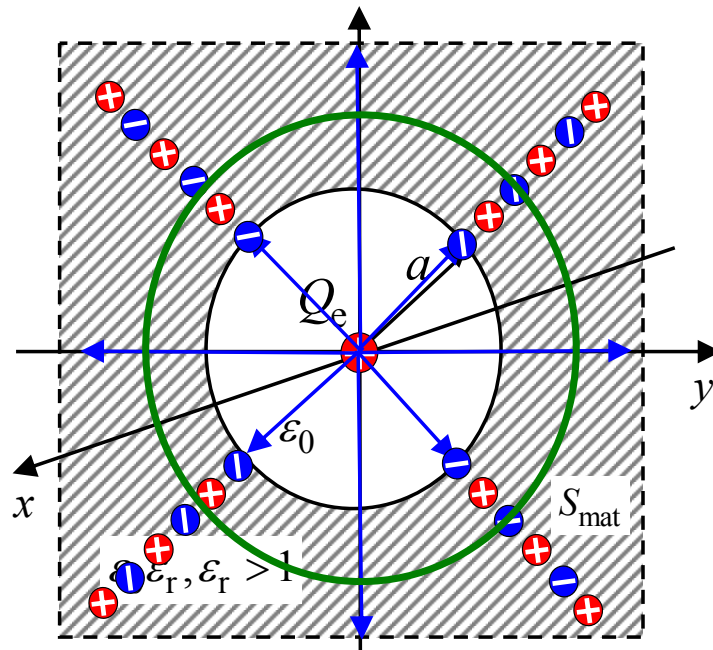
$$\underline{\mathbf{D}}(\underline{\mathbf{R}}) = \frac{Q_e}{4\pi} \frac{1}{R^2} \underline{\mathbf{e}}_R$$

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}) = \begin{cases} \frac{Q_e}{4\pi\varepsilon_0} \frac{1}{R^2} \underline{\mathbf{e}}_R & R < a \\ \frac{Q_e}{4\pi\varepsilon_0\varepsilon_r} \frac{1}{R^2} \underline{\mathbf{e}}_R & R > a \end{cases}$$

$$\underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = \begin{cases} \underline{\mathbf{0}} & R < a \\ \frac{Q_e}{4\pi} \frac{\varepsilon_r - 1}{\varepsilon_r} \frac{1}{R^2} \underline{\mathbf{e}}_R & R > a \end{cases}$$



ES Fields – Electric Polarization of Materials – Example / ES Felder – Elektrische Polarisierung von Materialien – Beispiel (...)



$$\oiint_{S_{mat}} \underline{\mathbf{D}}(\mathbf{R}) \cdot \underline{\mathbf{dS}} = Q_e$$

$$\underline{\mathbf{D}}(\mathbf{R}) = \epsilon_0 \underline{\mathbf{E}}(\mathbf{R}) + \underline{\mathbf{P}}_e(\mathbf{R})$$

$$\begin{aligned} \oiint_{S_{mat}} \underline{\mathbf{D}}(\mathbf{R}) \cdot \underline{\mathbf{dS}} &= \epsilon_0 \oiint_{S_{mat}} \underline{\mathbf{E}}(\mathbf{R}) \cdot \underline{\mathbf{dS}} + \oiint_{S_{mat}} \underline{\mathbf{P}}_e(\mathbf{R}) \cdot \underline{\mathbf{dS}} \\ &= Q_e \end{aligned}$$

$$\epsilon_0 \oiint_{S_{mat}} \underline{\mathbf{E}}(\mathbf{R}) \cdot \underline{\mathbf{dS}} = Q_e - \underbrace{\oiint_{S_{mat}} \underline{\mathbf{P}}_e(\mathbf{R}) \cdot \underline{\mathbf{dS}}}_{=-Q_e^{pol}}$$

$$= Q_e + Q_e^{pol}$$

$$Q_e = Q_e^{unpaired}$$

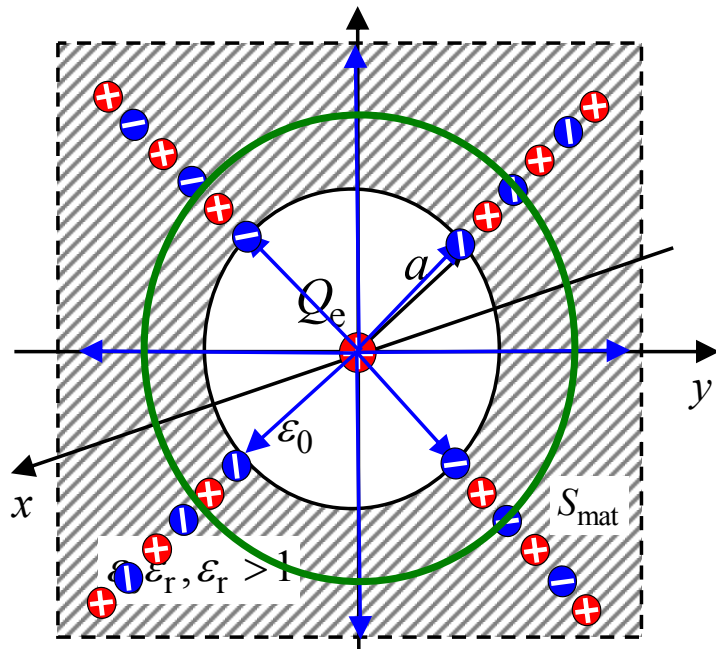
$$Q_e^{pol} = Q_e^{paired}$$



Unpaired Electric Charge /
Ungepaarte elektrische Ladungen

Paired Electric Charge /
Gepaarte elektrische Ladungen

ES Fields – Electric Polarization of Materials – Example / ES Felder – Elektrische Polarisierung von Materialien – Beispiel (...)



$$Q_e = \oiint_{S_{mat}} \underline{\mathbf{D}}(\underline{\mathbf{R}}) \cdot d\underline{\mathbf{S}}$$

$$= \iiint_{V_{mat}} \nabla \cdot \underline{\mathbf{D}}(\underline{\mathbf{R}}) dV$$

$$\nabla \cdot \underline{\mathbf{D}}(\underline{\mathbf{R}}) = \rho_e(\underline{\mathbf{R}})$$

$$= \rho_e^{\text{unpaired}}(\underline{\mathbf{R}})$$

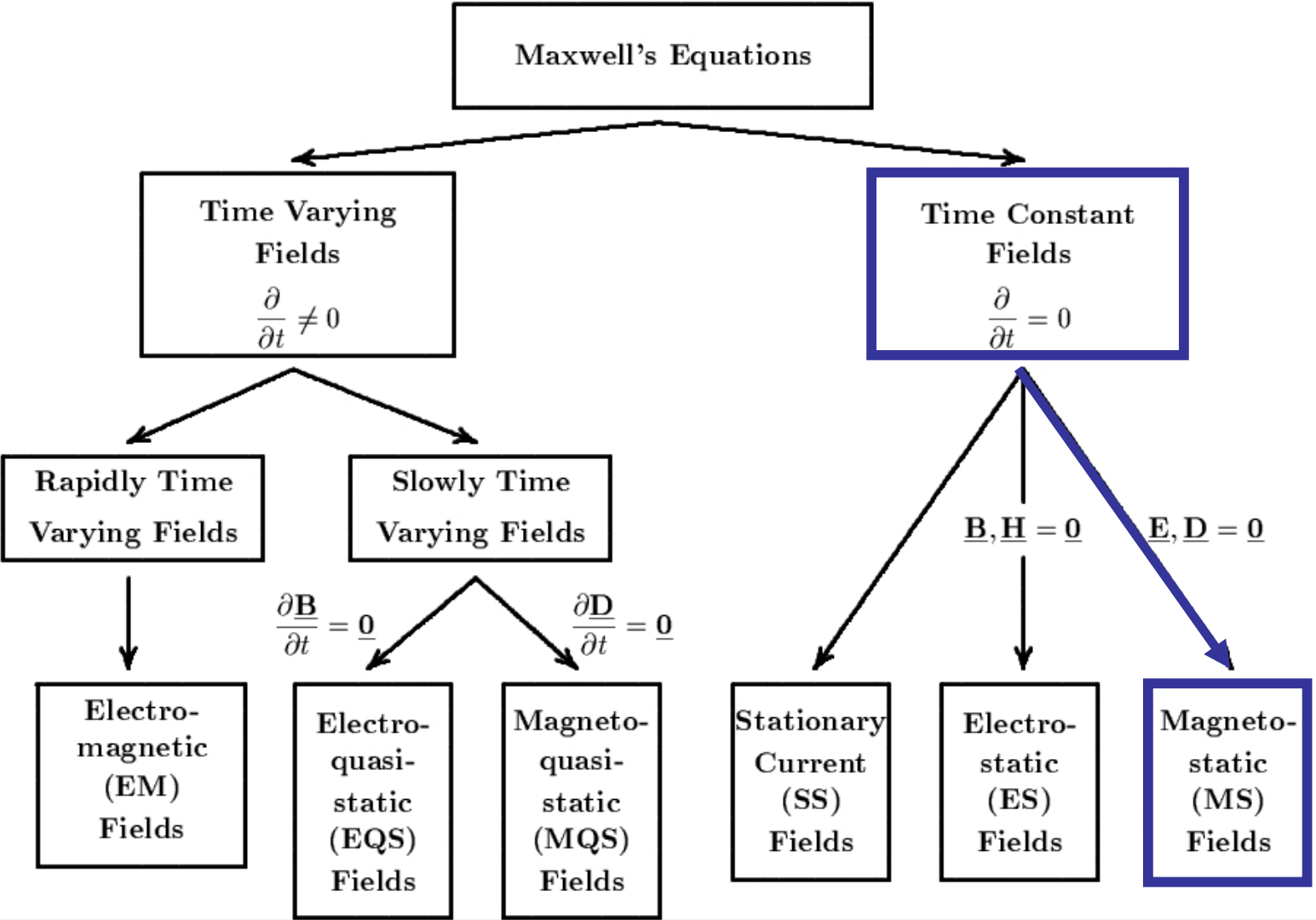
$$Q_e^{\text{pol}} = -\oiint_{S_{mat}} \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) \cdot d\underline{\mathbf{S}}$$

$$= -\iiint_{V_{mat}} \nabla \cdot \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) dV$$

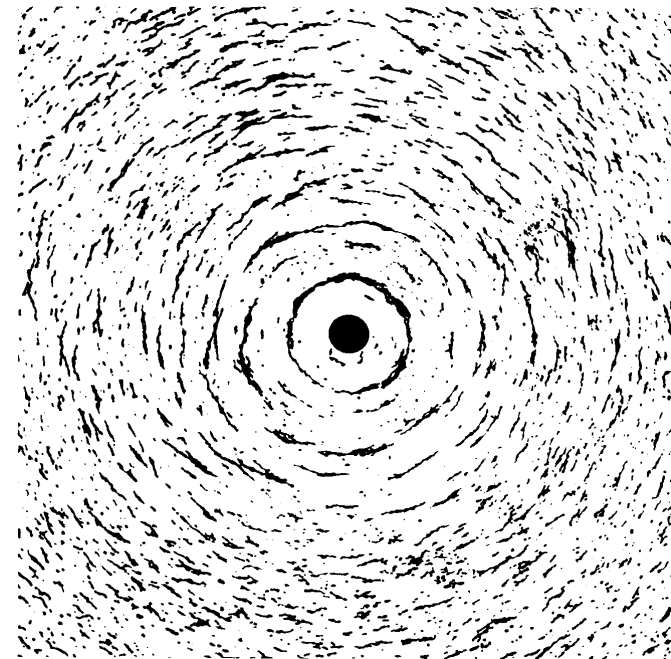
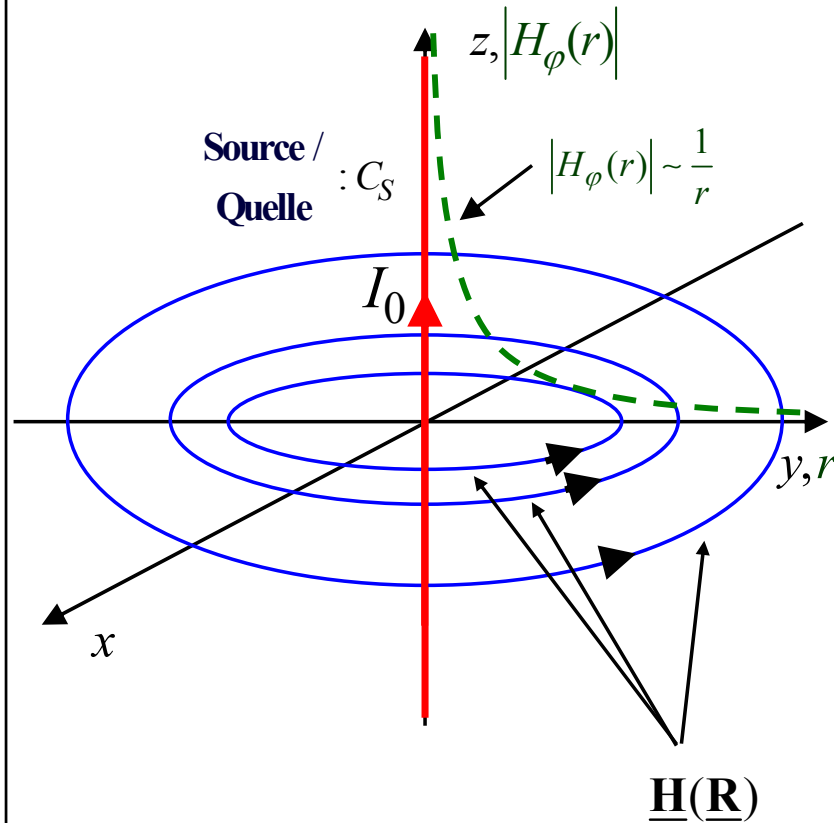
$$\nabla \cdot \underline{\mathbf{P}}_e(\underline{\mathbf{R}}) = -\rho_e^{\text{pol}}(\underline{\mathbf{R}})$$

$$= -\rho_e^{\text{paired}}(\underline{\mathbf{R}})$$

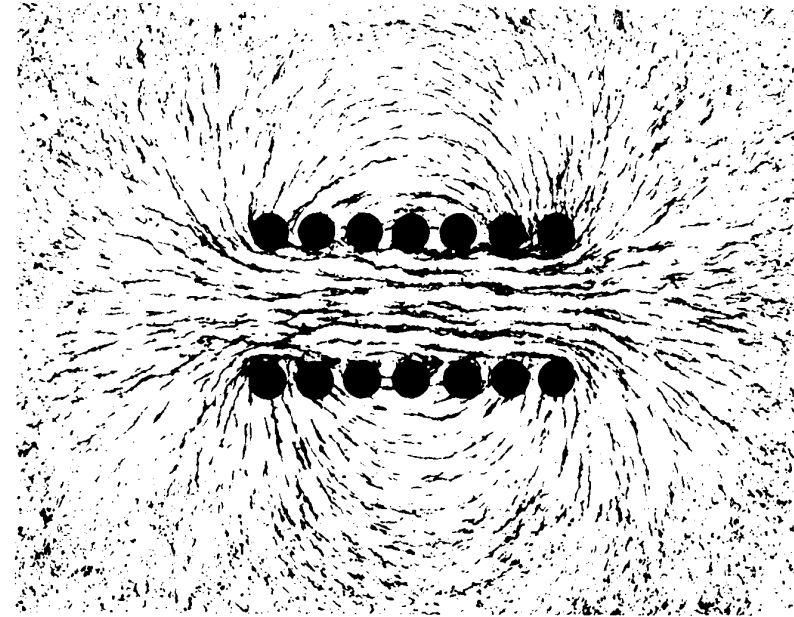
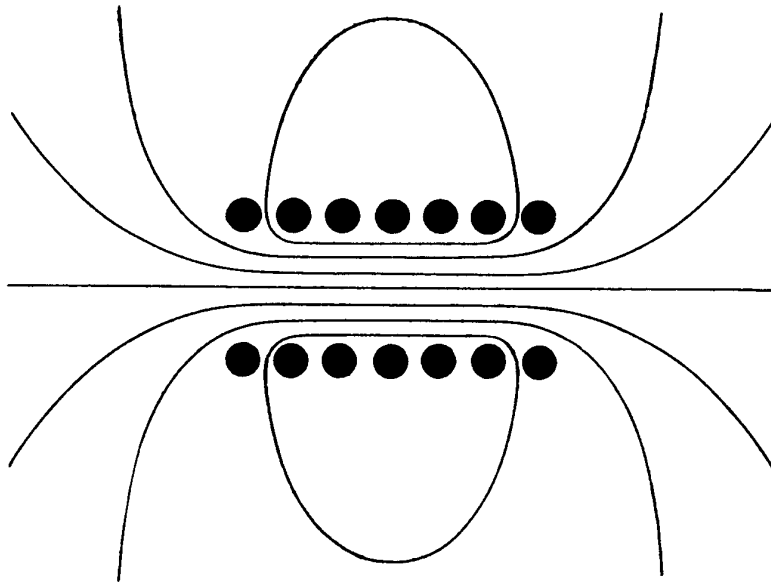
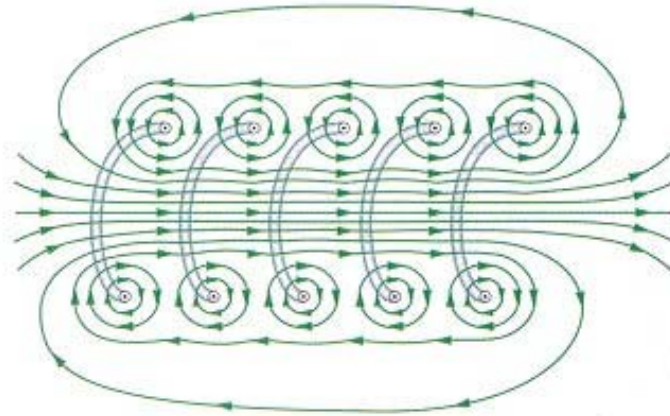
Magnetostatic (MS) Fields – Classification / Magnetostatische (MS) Felder - Klassifikation



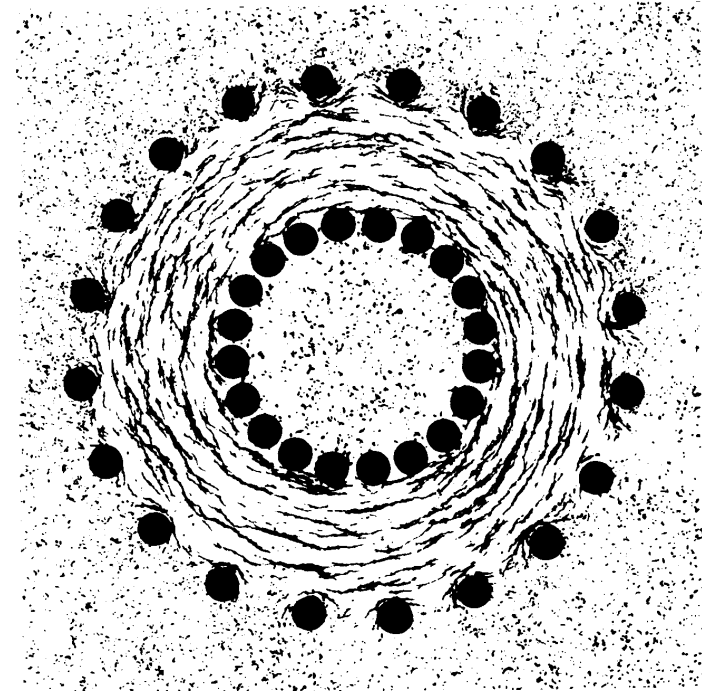
MS Fields – Magnetic Field of a Line Current / MS-Felder – Magnetfeld eines Linienstromes



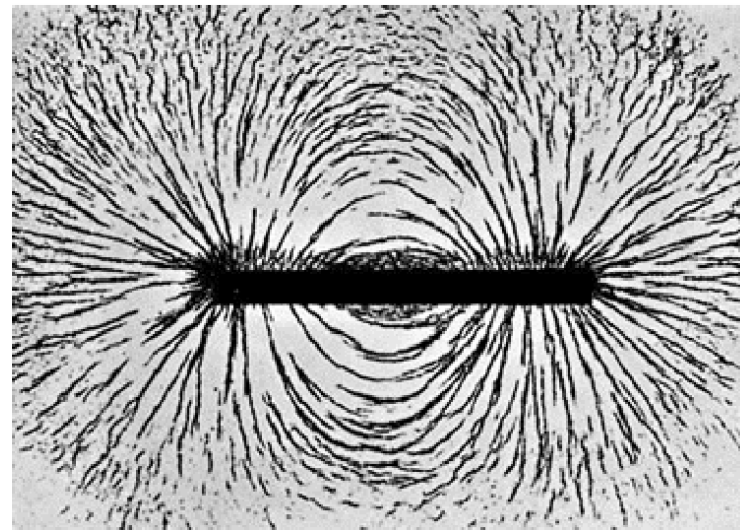
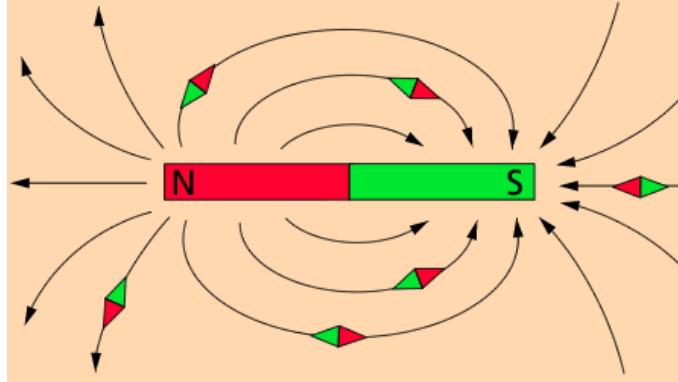
MS Fields – Magnetic Field of a Solenoid / MS-Felder – Magnetfeld einer Spule



MS Fields – Magnetic Field of a Toroidal Coil / MS-Felder – Magnetisches Feld einer Ringspule

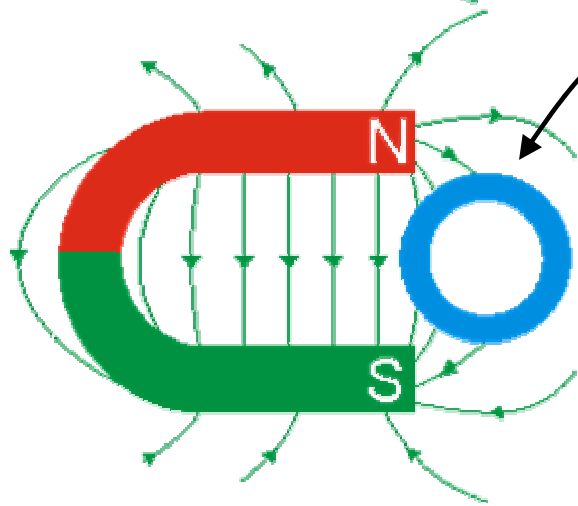
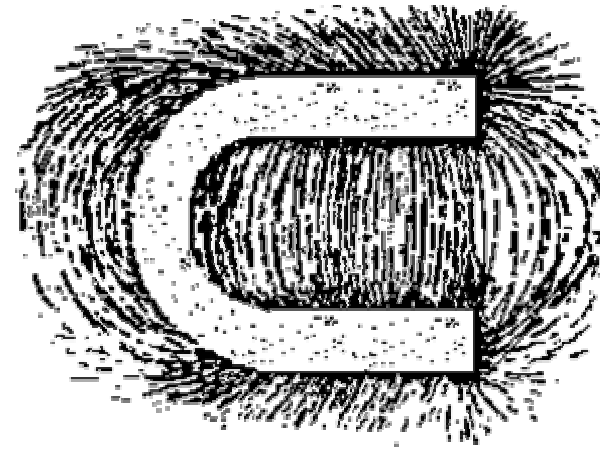
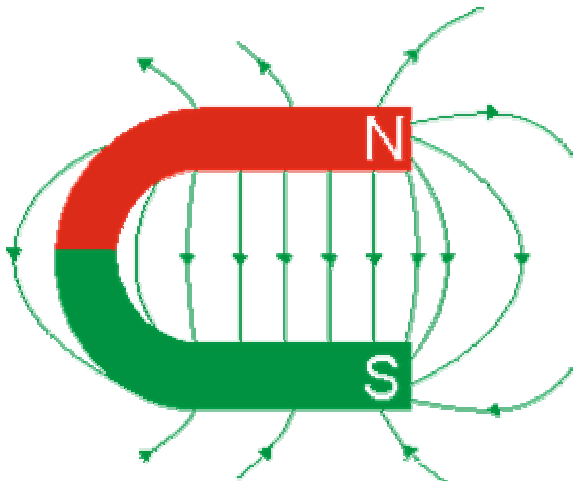


MS Fields – Magnetic Field of a Bar Magnet / MS-Felder – Magnetfeld eines Stabmagneten

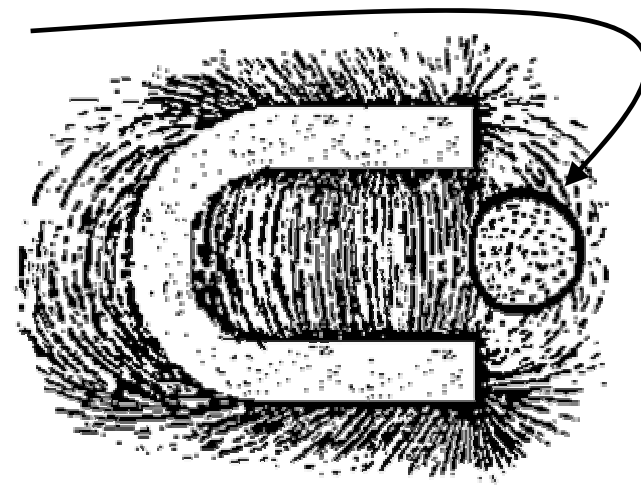


**Magnetic Field Lines of a Bar Magnet – Visualized with Iron Filings /
Magnetfeldlinien eines Stabmagneten – Visualisiert mit Eisenfeilspäne**

**MS Fields – Magnetic Fields of a Horseshoe Magnet (Permanent Magnet) /
MS-Felder – Magnetfeld eines Hufeisenmagneten (Permanentmagnet)**

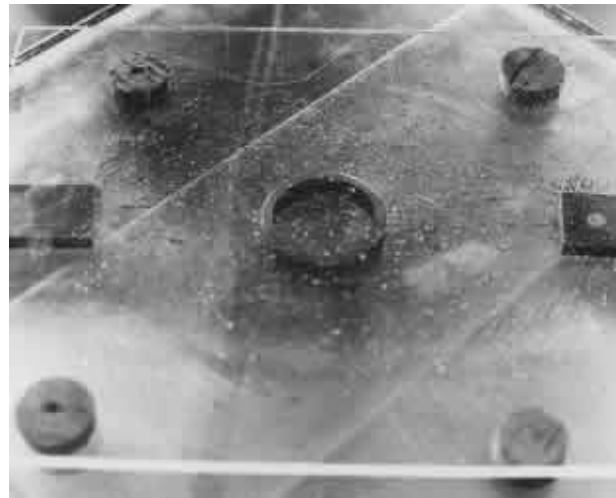
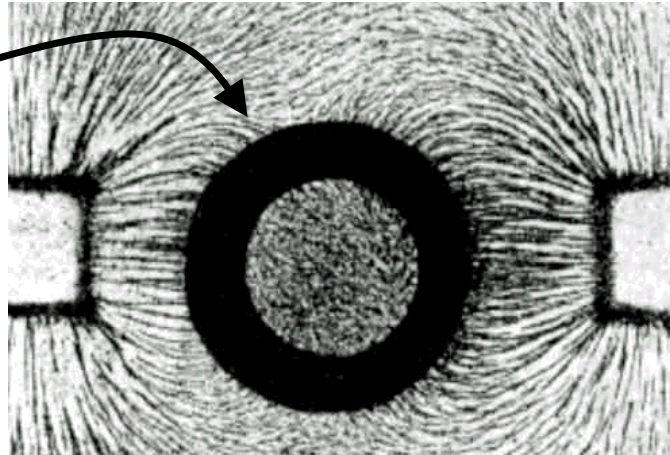


Iron Ring /
Eiserring

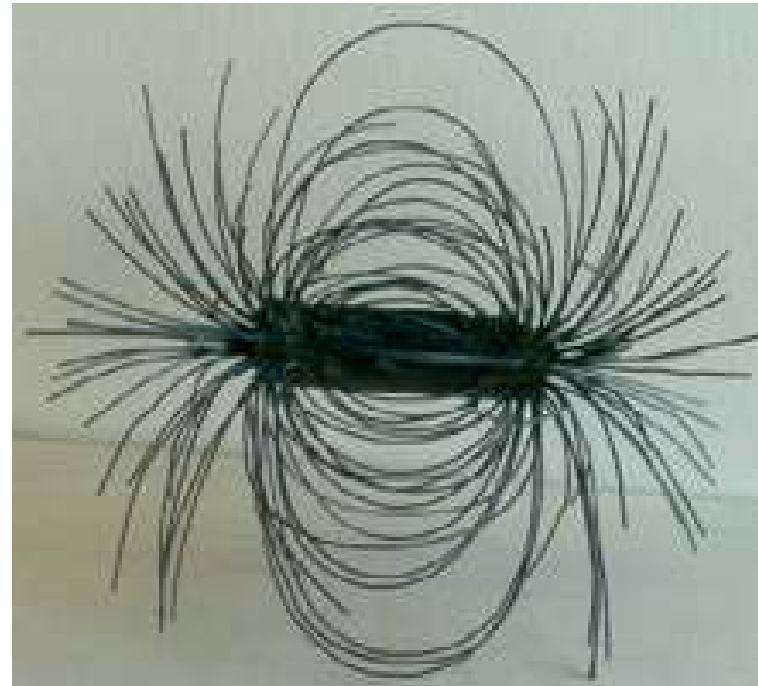


MS Fields – (...) Magnetic Shielding / MS-Felder – (...) Magnetische Abschirmung

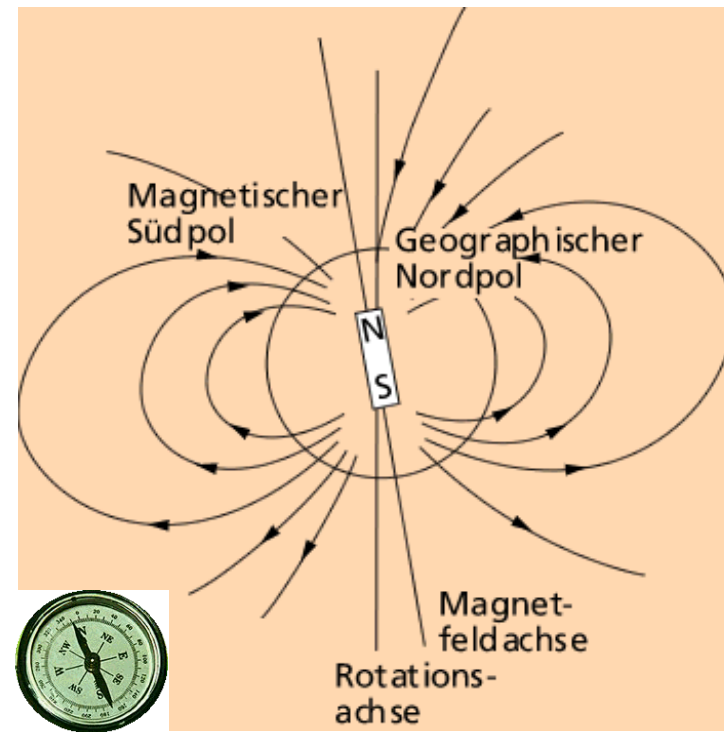
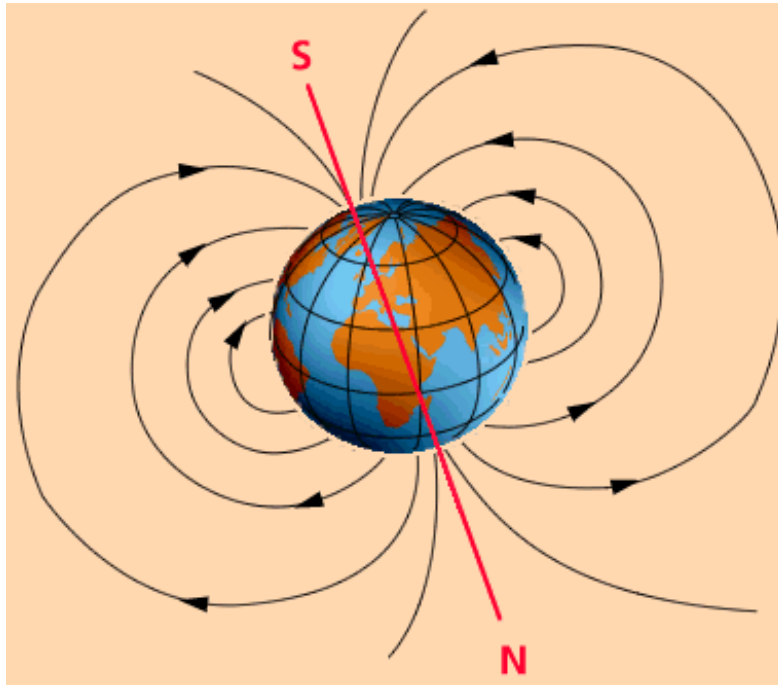
Iron Ring /
Eisenring



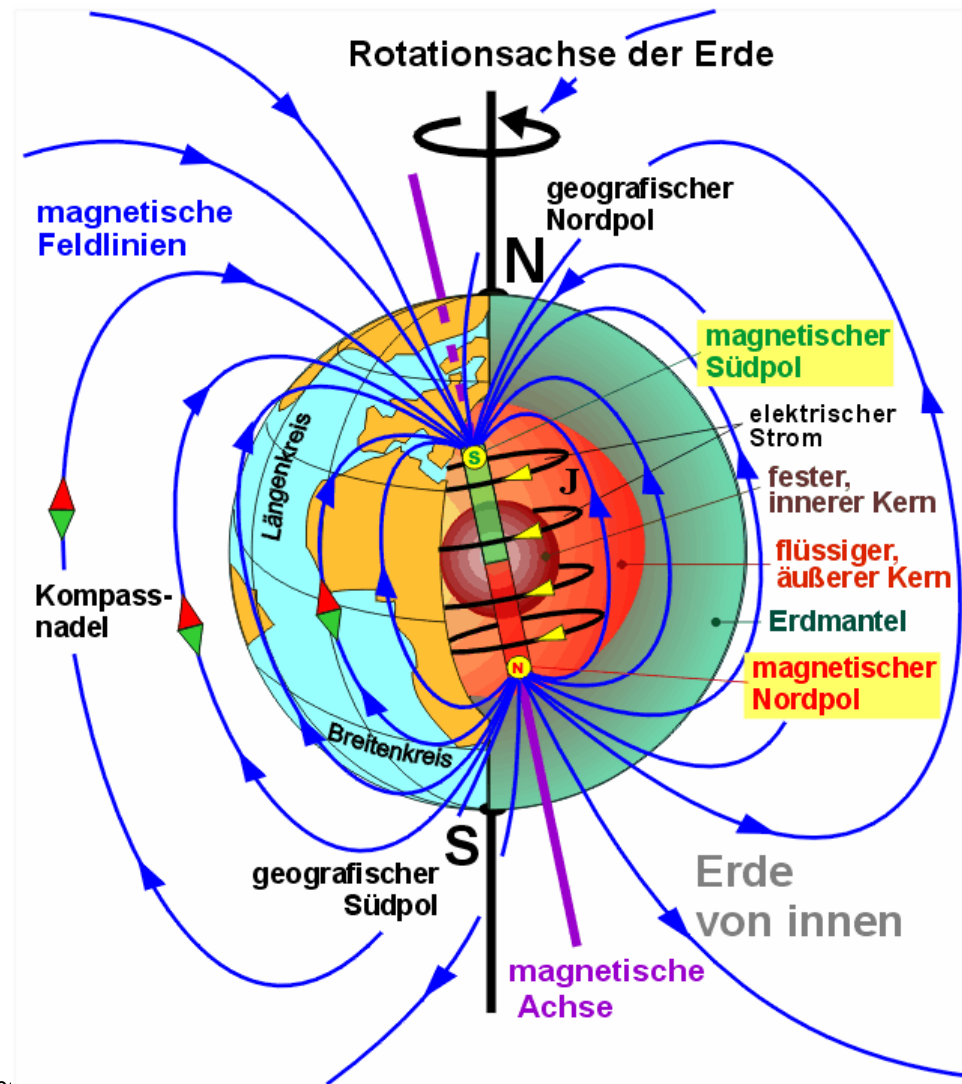
MS Fields – Magnetic Fields of Permanent Magnets / MS-Felder – Magnetfelder von Permanentmagneten



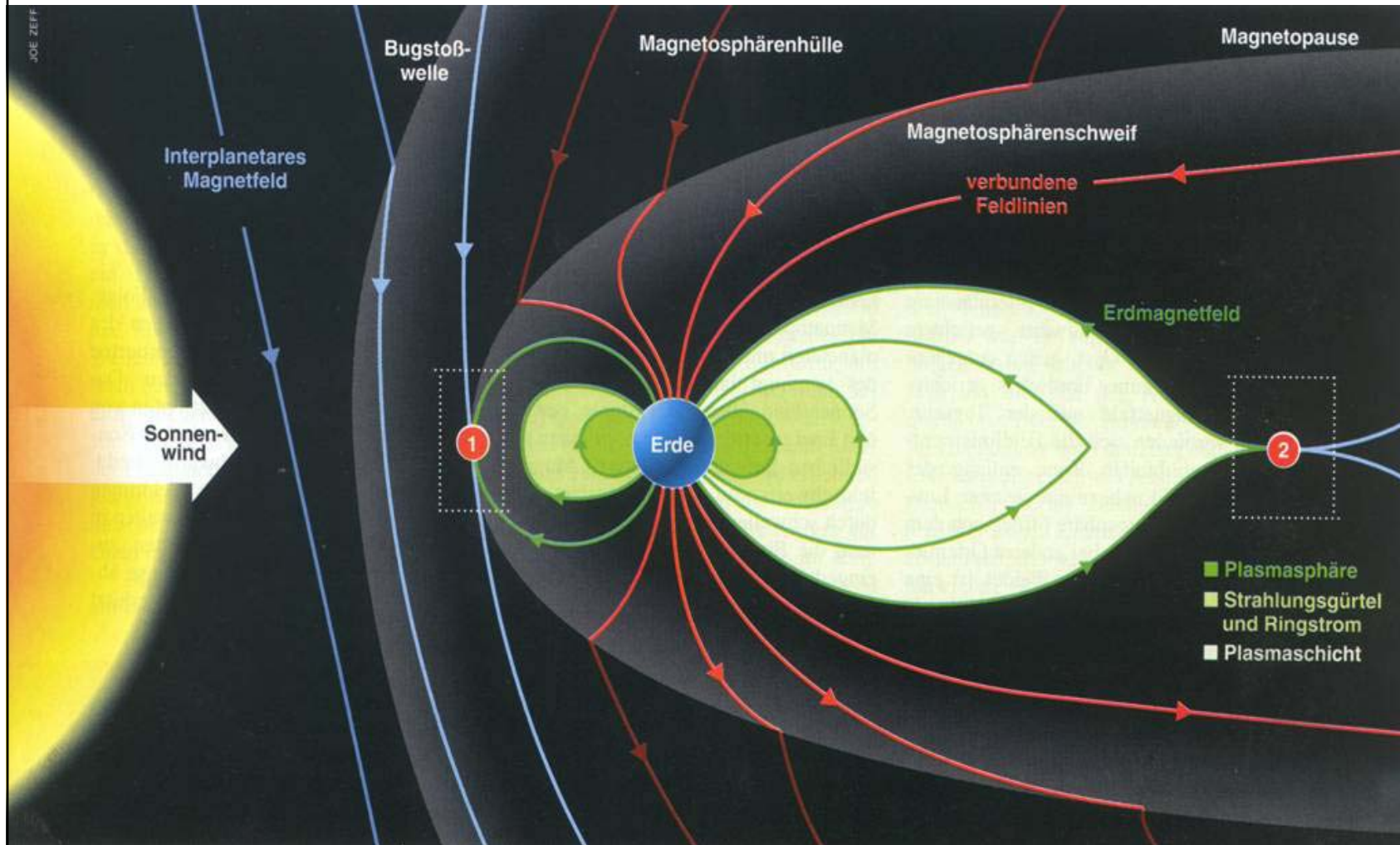
MS Fields – Magnetic Field of the Earth / MS-Felder – Magnetfeld der Erde



MS Fields – Magnetic Field of the Earth / MS-Felder – Magnetfeld der Erde



MS Fields – Magnetic Field of the Earth – Sun Wind / MS-Felder - Magnetfeld der Erde - Sonnenwind



MS Fields – Governing Equations / MS-Felder - Grundgleichungen

**Integral Form /
Integralform**

$$\oint_{C=\partial S} \underline{\mathbf{H}}(\underline{\mathbf{R}}) \cdot d\underline{\mathbf{R}} = \iint_S \underline{\mathbf{J}}_e(\underline{\mathbf{R}}) \cdot d\underline{\mathbf{S}}$$

$$\oiint_S \underline{\mathbf{B}}(\underline{\mathbf{R}}) \cdot d\underline{\mathbf{S}} = 0$$

**Differential Form /
Differentialform**

$$\nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}) = \underline{\mathbf{J}}_e(\underline{\mathbf{R}})$$

$$\nabla \cdot \underline{\mathbf{B}}(\underline{\mathbf{R}}) = 0$$

$\underline{\mathbf{J}}_e(\underline{\mathbf{R}})$

is a Known Prescribed Electric Current Density: For Example a Electric Current Density in a Wire /
ist eine bekannte vorgegebene elektrische Stromdichte: Zum Beispiel eine elektrische Stromdichte in einem Draht

**Because of /
Weil**

$$\nabla \cdot \underline{\mathbf{B}}(\underline{\mathbf{R}}) = 0$$

**$\underline{\mathbf{B}}(\underline{\mathbf{R}})$ Can be Represented by /
kann dargestellt werden über $\underline{\mathbf{B}}(\underline{\mathbf{R}}) = \nabla \times \underline{\mathbf{A}}_e(\underline{\mathbf{R}})$**

**In Vacuum we have /
Im Vakuum gilt**

$$\underline{\mathbf{B}}(\underline{\mathbf{R}}) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}})$$

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{1}{\mu_0} \underline{\mathbf{B}}(\underline{\mathbf{R}})$$

$$= \frac{1}{\mu_0} \nabla \times \underline{\mathbf{A}}_e(\underline{\mathbf{R}})$$

MS Fields – Vector Potential / MS Felder – Vektorpotential

$$\nabla \cdot \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) = 0 \quad \text{Coulomb Gauge / Coulomb-Eichung}$$

It follows /
Es folgt

$$\mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}}) = \nabla \times \underline{\mathbf{A}}_e(\underline{\mathbf{R}})$$

Applying the Curl Operator Gives /
Die Anwendung des Rotationsoperators ergibt

$$\begin{aligned} \mu_0 \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}) &= \nabla \times \nabla \times \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) \\ &= \mu_0 \underline{\mathbf{J}}_e(\underline{\mathbf{R}}) \end{aligned}$$

$$\begin{aligned} \nabla \times \nabla \times \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) &= \nabla \underbrace{\nabla \cdot \underline{\mathbf{A}}_e(\underline{\mathbf{R}})}_{=0} - \nabla \cdot \nabla \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) \\ &= -\nabla \cdot \nabla \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) \\ &= -\Delta \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) \end{aligned}$$

MS Fields – Vector Potential – Poisson and Laplace Equation / MS Felder – Vektorpotential – Poisson- und Laplace-Gleichung

$$\Delta \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) = \begin{cases} -\mu_0 \underline{\mathbf{J}}_e(\underline{\mathbf{R}}) & \text{for / für } \underline{\mathbf{J}}_e(\underline{\mathbf{R}}) \neq \underline{\mathbf{0}} & \text{Vectorial Poisson Equation /} \\ & & \text{Vektorielle Poisson-Gleichung} \\ 0 & \text{for / für } \underline{\mathbf{J}}_e(\underline{\mathbf{R}}) = \underline{\mathbf{0}} & \text{Vectorial Laplace Equation /} \\ & & \text{Vektorielle Laplace-Gleichung} \end{cases}$$

MS Fields – Vector Potential – (...) Special Solution / MS Felder – Vector Potential – (...) Spezielle Lösung

$$\Delta \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) = -\mu_0 \underline{\mathbf{J}}_e(\underline{\mathbf{R}}) \quad \text{Vector Poisson's Equation /} \\ \text{Vektorielle Poisson-Gleichung}$$

Special Solution /
Spezielle Lösung



$$\underline{\mathbf{A}}_e(\underline{\mathbf{R}}) = \mu_0 \iiint_{V_S} G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') \underline{\mathbf{J}}_e(\underline{\mathbf{R}}') d^3 \underline{\mathbf{R}}'$$

With the Three-Dimensional
Static Green's Function /
Mit der dreidimensionalen
statischen Greenschen Funktion

$$G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') = \frac{1}{4\pi} \frac{1}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|}$$

$$\begin{aligned} \underline{\mathbf{H}}(\underline{\mathbf{R}}) &= \frac{1}{\mu_0} \nabla \times \underline{\mathbf{A}}_e(\underline{\mathbf{R}}) \\ &= \frac{1}{\mu_0} \nabla \times \left[\mu_0 \iiint_{V_S} G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') \underline{\mathbf{J}}_e(\underline{\mathbf{R}}') d^3 \underline{\mathbf{R}}' \right] \\ &= \iiint_{V_S} \nabla \times \left[G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') \underline{\mathbf{J}}_e(\underline{\mathbf{R}}') \right] d^3 \underline{\mathbf{R}}' \end{aligned}$$

$$\nabla \times (\Phi \underline{\mathbf{A}}) = \Phi \nabla \times \underline{\mathbf{A}} + \nabla \Phi \times \underline{\mathbf{A}}$$

$$\nabla \times [G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') \underline{\mathbf{J}}_e(\underline{\mathbf{R}}')] = G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') \underbrace{\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}')}_{=0} + \nabla G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}')$$

MS Fields – Vector Potential – Biot-Savart's Law / MS Felder – Vector Potential – Biot-Savartsche Gesetz

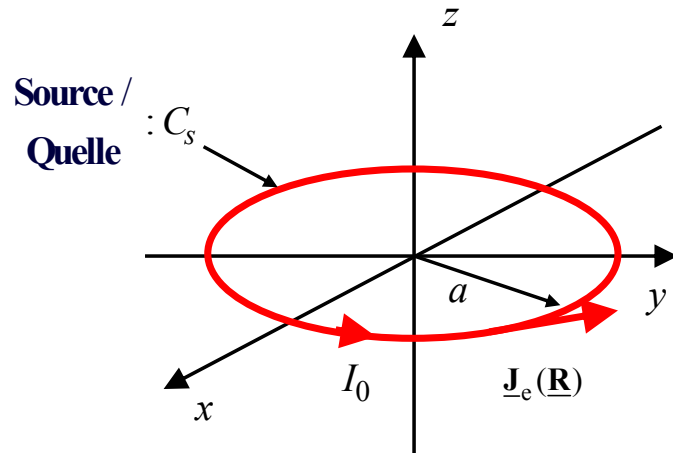
$$\nabla G(\underline{\mathbf{R}} - \underline{\mathbf{R}}') = \nabla \frac{1}{4\pi} \frac{1}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|} = \frac{1}{4\pi} \nabla \frac{1}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|} = -\frac{1}{4\pi} \frac{\underline{\mathbf{R}} - \underline{\mathbf{R}}'}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3}$$

$$\begin{aligned} \underline{\mathbf{H}}(\underline{\mathbf{R}}) &= -\iiint_{V_S} \frac{1}{4\pi} \frac{\underline{\mathbf{R}} - \underline{\mathbf{R}}'}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}') d^3 \underline{\mathbf{R}}' \\ &= \iiint_{V_S} \frac{1}{4\pi} \frac{\underline{\mathbf{J}}_e(\underline{\mathbf{R}}') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} d^3 \underline{\mathbf{R}}' \\ &= \frac{1}{4\pi} \iiint_{V_S} \frac{\underline{\mathbf{J}}_e(\underline{\mathbf{R}}') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} d^3 \underline{\mathbf{R}}' \end{aligned}$$

Biot-Savart's Law (for a Given Volume Source) /
Biot-Savartsches Gesetz (für eine Volumenquelle)

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{1}{4\pi} \iiint_{V_S} \frac{\underline{\mathbf{J}}_e(\underline{\mathbf{R}}') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} d^3 \underline{\mathbf{R}}'$$

MS Fields – Biot-Savart's Law for a Line Source / MS Felder – Biot-Savartsches Gesetz für eine Linienquelle



$$\underline{\mathbf{J}}_e(\underline{\mathbf{R}}) = I_0 \delta(r - a) \delta(z) \underline{\mathbf{e}}_{\varphi}, \quad a > 0$$

$$\begin{aligned} \underline{\mathbf{H}}(\underline{\mathbf{R}}) &= \frac{1}{4\pi} \iiint_{V_s} \frac{\underline{\mathbf{J}}_e(\underline{\mathbf{R}}') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} d^3 \underline{\mathbf{R}}' \\ &= \frac{1}{4\pi} \int_{z'=-\infty}^{\infty} \int_{\varphi'=0}^{2\pi} \int_{r'=0}^{\infty} \frac{I_0 \delta(r' - a) \delta(z') \underline{\mathbf{e}}_{\varphi'}(\varphi') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} r' dr' d\varphi' dz' \\ &= \frac{I_0}{4\pi} \int_{\varphi'=0}^{2\pi} \int_{r'=0}^{\infty} \frac{\delta(r' - a) \underline{\mathbf{e}}_{\varphi'}(\varphi') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} r' dr' d\varphi' \Big|_{\underline{\mathbf{R}}': z'=0} \end{aligned}$$

MS Fields – Biot-Savart's Law for a Line Source / MS-Felder – Biot-Savartsche Gesetz für eine Linienquelle

$$\begin{aligned}
 \underline{\mathbf{H}}(\underline{\mathbf{R}}) &= \frac{I_0}{4\pi} \int_{\varphi'=0}^{2\pi} \frac{\overbrace{\underline{\mathbf{e}}_{\varphi'}(\varphi') a d\varphi'}^{=d\underline{\mathbf{R}}'} \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} \Bigg|_{\underline{\mathbf{R}}': r'=a; z'=0} \\
 &= \frac{I_0}{4\pi} \int_{\varphi'=0}^{2\pi} \frac{d\underline{\mathbf{R}}' \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} \Bigg|_{\underline{\mathbf{R}}': r'=a; z'=0} \\
 &= \frac{I_0}{4\pi} \int_{C_S} \frac{d\underline{\mathbf{R}}' \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} \Bigg|_{C_S: \underline{\mathbf{R}}': r'=a; 0 \leq \varphi' < 2\pi; z'=0}
 \end{aligned}$$

Line Source with Contour C_S /
Linienquelle mit der Kontur C_S $C_S : \underline{\mathbf{R}}' : r' = a; 0 \leq \varphi' < 2\pi; z' = 0$

Biot-Savart Law for a Line Source with Contour C_S /
Biot-Savartsches Gesetz für eine Linienquelle mit der Kontur C_S

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{I_0}{4\pi} \int_{C_S} \frac{d\underline{\mathbf{R}}' \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3}$$



Arbitrary Contour C_S /
Beliebige Kontur C_S

MS Fields – Biot-Savart's Law for a Line Source / MS-Felder – Biot-Savartsche Gesetz für eine Linienquelle

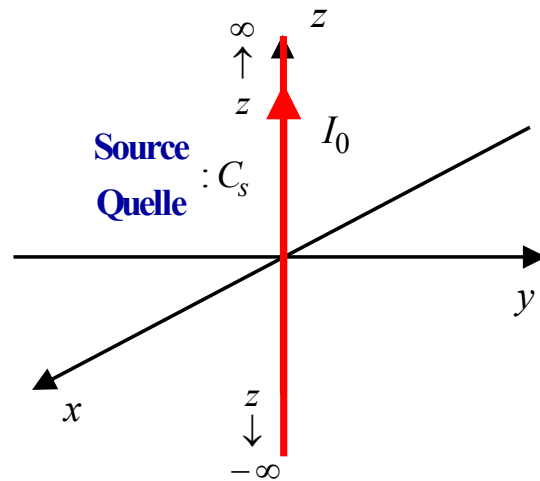
Biot-Savart's Law for a Volume Source V_s /
Biot-Savartsches Gesetz für eine Volumenquelle V_s

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{1}{4\pi} \iiint_{V_s} \frac{\underline{\mathbf{J}}_e(\underline{\mathbf{R}}') \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3} d^3 \underline{\mathbf{R}}'$$

Biot-Savart Law for a Line Source with Contour C_s /
Biot-Savartsches Gesetz für eine Linienquelle mit der Kontur C_s

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{I_0}{4\pi} \int_{C_s} \frac{d\underline{\mathbf{R}}' \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3}$$

MS Fields – Biot-Savart – Example: Infinite Thin and Infinite Long Wire Carrying a Constant Electric Current / MS-Felder – Biot-Savart – Beispiel: unendlich dünner, unendlich langer Draht, der einen konstanten elektrischen Strom führt



Biot-Savart Law for a Line Source with Contour C_S /
Biot-Savartsches Gesetz für eine Linienquelle mit der Kontur C_S

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{I_0}{4\pi} \int_{C_S} \frac{d\underline{\mathbf{R}}' \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}')}{|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3}$$

$$\underline{\mathbf{R}} = r\underline{\mathbf{e}}_r(\varphi) + z\underline{\mathbf{e}}_z$$

$$\begin{aligned} \underline{\mathbf{R}}' &= r'\underline{\mathbf{e}}_{r'}(\varphi') + z'\underline{\mathbf{e}}_z \Big|_{r'=0} \\ &= z'\underline{\mathbf{e}}_z, \quad -\infty \leq z' \leq \infty \end{aligned}$$

$$\underline{\mathbf{R}} - \underline{\mathbf{R}}' = r\underline{\mathbf{e}}_r + (z - z')\underline{\mathbf{e}}_z$$

$$\begin{aligned} d\underline{\mathbf{R}}' &= \frac{d}{dz'} \underline{\mathbf{R}}' dz' \\ &= \underline{\mathbf{e}}_z dz' \end{aligned}$$

$$\begin{aligned} d\underline{\mathbf{R}}' \times (\underline{\mathbf{R}} - \underline{\mathbf{R}}') &= \begin{vmatrix} \underline{\mathbf{e}}_r & \underline{\mathbf{e}}_\varphi(\varphi) & \underline{\mathbf{e}}_z \\ 0 & 0 & dz' \\ r & 0 & z - z' \end{vmatrix} \\ &= r \underline{\mathbf{e}}_\varphi dz' \end{aligned}$$

$$|\underline{\mathbf{R}} - \underline{\mathbf{R}}'|^3 = [r^2 + (z - z')^2]^{3/2}$$

MS Fields – Biot-Savart – Example: Infinite Thin and Infinite Long Wire Carrying a Constant Electric Current / MS-Felder – Biot-Savart – Beispiel: unendlich dünn, unendlich langer Draht, der einen konstanten elektrischen Strom führt (...)

$$\begin{aligned}
 \underline{\mathbf{H}}(\underline{\mathbf{R}}) &= \frac{I_0}{4\pi} \int_{z'=-\infty}^{\infty} \frac{r \underline{\mathbf{e}}_{\varphi}(\varphi)}{\left[r^2 + (z - z')^2\right]^{3/2}} dz' \\
 &= -\frac{I_0}{4\pi} r \underline{\mathbf{e}}_{\varphi}(\varphi) \int_{\alpha=\infty}^{-\infty} \frac{d\alpha}{\left[r^2 + \alpha^2\right]^{3/2}} \\
 &= \frac{I_0}{4\pi} r \underline{\mathbf{e}}_{\varphi}(\varphi) \int_{\alpha=-\infty}^{\infty} \frac{d\alpha}{\left[r^2 + \alpha^2\right]^{3/2}} \\
 &= \frac{I_0}{4\pi} r \underline{\mathbf{e}}_{\varphi}(\varphi) \int_{\alpha=-\infty}^{\infty} \frac{d\alpha}{\left[r^2 + \alpha^2\right]^{3/2}} \\
 &= \frac{I_0}{4\pi} r \underline{\mathbf{e}}_{\varphi}(\varphi) \frac{\alpha}{r^2 \sqrt{r^2 + \alpha^2}} \Bigg|_{\alpha=-\infty}^{\infty}
 \end{aligned}$$

With the Substitution / Mit der Substitution

$$\begin{aligned}
 \alpha &= z - z' \\
 dz' &= -d\alpha \\
 z' = -\infty &: \quad \alpha = z + \infty \\
 &= \infty \\
 z' = \infty &: \quad \alpha = z - \infty \\
 &= -\infty
 \end{aligned}$$

$$\int \frac{dx}{\left[ax^2 + b\right]^{3/2}} = \frac{x}{b\sqrt{ax^2 + b}}$$

with
mit $a = 1, b = r^2$

MS Fields – Biot-Savart – Example: Infinite Thin and Infinite Long Wire Carrying a Constant Electric Current / MS-Felder – Biot-Savart – Beispiel: unendlich dünner, unendlich langer Draht, der einen konstanten elektrischen Strom führt (...)

$$\underline{\mathbf{H}}(\mathbf{R}) = \frac{I_0}{4\pi r} \frac{\text{sgn}(\alpha) |\alpha|}{\sqrt{r^2 + \alpha^2}} \Bigg|_{\alpha=-\infty}^{\infty} \underline{\mathbf{e}}_{\varphi}(\varphi)$$

$$= \frac{I_0}{4\pi r} \frac{\text{sgn}(\alpha) |\alpha|}{|\alpha| \sqrt{1 + r^2 / \alpha^2}} \Bigg|_{\alpha=-\infty}^{\infty} \underline{\mathbf{e}}_{\varphi}(\varphi)$$

$$= \frac{I_0}{4\pi r} \left[\underbrace{\lim_{\alpha \rightarrow \infty} \frac{\text{sgn}(\alpha)}{\sqrt{1 + r^2 / \alpha^2}}}_{=1} - \underbrace{\lim_{\alpha \rightarrow -\infty} \frac{\text{sgn}(\alpha)}{\sqrt{1 + r^2 / \alpha^2}}}_{=-1} \right] \underline{\mathbf{e}}_{\varphi}(\varphi)$$

$$= \frac{I_0}{4\pi r} \underbrace{[1 - (-1)]}_2 \underline{\mathbf{e}}_{\varphi}(\varphi)$$

$$= \frac{I_0}{2\pi r} \underline{\mathbf{e}}_{\varphi}(\varphi)$$

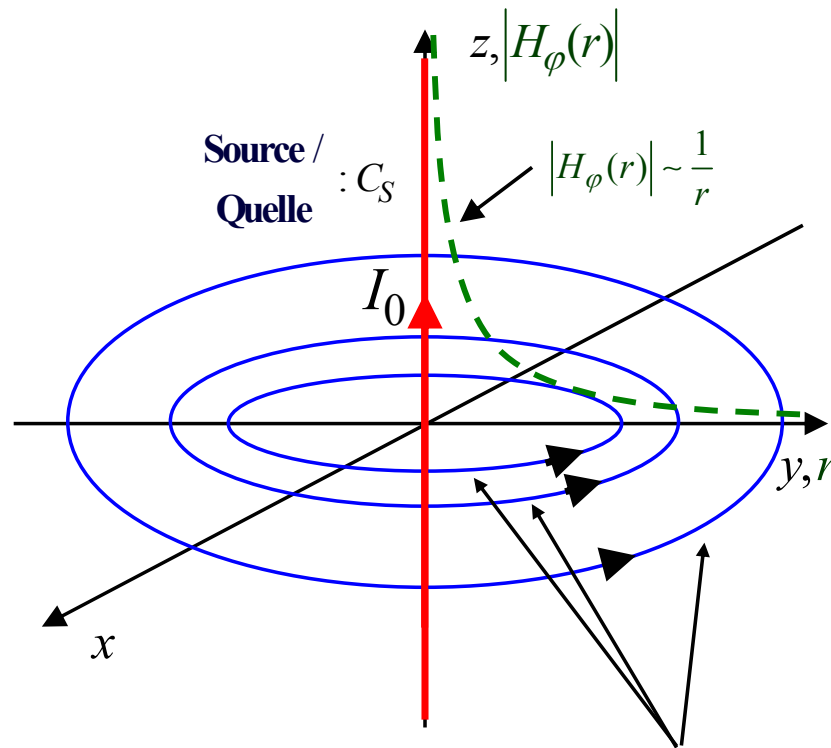
With the Signum Function /
Mit der Signum-Funktion

$$\text{sgn}(\alpha) = \begin{cases} -1 & \alpha < 0 \\ 1 & \alpha > 0 \end{cases}$$

$$\alpha = \text{sgn}(\alpha) |\alpha|$$

With /
Mit $\sqrt{\alpha^2} = |\alpha|$

MS Fields – Biot-Savart – Example: Infinite Thin and Infinite Long Wire Carrying a Constant Electric Current / MS-Felder – Biot-Savart – Beispiel: unendlich dünner, unendlich langer Draht, der einen konstanten elektrischen Strom führt (...)



$$\underline{\mathbf{H}}(\underline{\mathbf{R}}) = \frac{I_0}{2\pi r} \underline{\mathbf{e}}_\varphi(\varphi)$$

$$= H_\varphi(r) \underline{\mathbf{e}}_\varphi(\varphi)$$

$$|\underline{\mathbf{H}}(\underline{\mathbf{R}})| = |H_\varphi(r)| \sim \frac{1}{r}$$

Field Lines of the Magnetic Field Strength / Feldlinien der magnetischen Feldstärke $\underline{\mathbf{H}}(\underline{\mathbf{R}})$

End of the 10th Lecture / Ende der 10. Vorlesung